

**“Analysis of 4H Silicon Carbide Double Implanted MOSFET  
using Distorted Gaussian Doping Profile in Drift Region for  
Enhanced Breakdown Voltage and Low Power Dissipation”**

*A Dissertation submitted in partial fulfillment towards the  
requirement for the degree of*

**MASTER OF TECHNOLOGY  
IN  
VLSI DESIGN**

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DEPARTMENT OF ELECTRONICS AND COMMUNICATION ENGINEERING


**DECLARATION**

I hereby declare that the work which is being presented in the dissertation titled “**Analysis of 4H Silicon Carbide Double Implanted MOSFET using Distorted Gaussian Doping Profile in Drift Region for Enhanced Breakdown Voltage and Low Power Dissipation**” in partial fulfillment of the requirement for the award of degree of Master of Technology in VLSI Design submitted in Electronics and Communication Engineering Department of Thapar University, Patiala is an authentic record of my study carried out under the guidance of **Dr.A.K.Chatterjee**(Professor, ECED) during 2014-2016.

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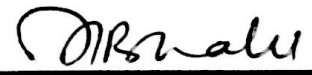
It is certified that the above statement made by the student is correct to the best of my knowledge and belief.

  
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Nickson P Jose

# Abstract

Silicon Carbide, unlike Silicon, is a Wide Band Gap (WBG) semiconductor that carries with it many properties like high thermal conductivity, high electric field strength and high saturation drift velocity, making it the ideal material for electronic industry. The ability of Silicon Carbide, like Silicon, to easily form high quality native oxide( $\text{SiO}_2$ ), gives it an advantage over other WBG semiconductors like Gallium Nitride and Aluminium Nitride. The research works related to SiC started in late 1980s but received a setback owing to the absence of technology to extract single crystal substrates. However, the above issue and also the lack of reproducible techniques has been taken care of by rapidly maturing extraction technology.

The work carried out here on 4H-Silicon Carbide Double Implanted MOSFET(DIMOSFET) is to analyze the performance of the above device related to power dissipation and breakdown voltage for distorted Gaussian doping profile in the drift region. The doping profile which is used here is a more practical profile (that can be obtained by mechanism like Molecular Beam Epitaxy(MBE)). This profile helps to increase the breakdown voltage while at the same time reduces the series parasitic resistance at the lower end of the device thereby lowering the power dissipation, all of which has also been successfully observed in theoretical analysis here.

# List of Figures

1.1	Stacking sequence of the most common polytypes of SiC . . . . .	3
1.2	Structure of DIMOSFET . . . . .	6
1.3	Resistance distribution in DIMOSFET . . . . .	7
1.4	Cross sectional view of DIMOSFET . . . . .	7
3.1	Uniformly Doped Profile in DIMOSFET . . . . .	19
3.2	Linearly Graded Doping Profile in DIMOSFET . . . . .	20
4.1	Distorted Gaussian Doping Profile in DIMOSFET . . . . .	22
4.2	Basic Device Structure of DIMOSFET . . . . .	23
4.3	Distorted Gaussian Profile . . . . .	24
4.4	Plot of current density vs power dissipation for 4H-SiC DIMOSFET for different values of $N_o$ . . . . .	30
4.5	Plot of current density vs on-state depletion width for 4H-SiC DIMOSFET for different values of $N_o$ . . . . .	30
4.6	Plot of current density vs drain-to-source voltage for 4H-SiC DIMOSFET for different values of $N_o$ . . . . .	31
4.7	Plot of effective doping concentration vs concentration gradient(region I) for 4H-SiC DIMOSFET . . . . .	31
4.8	Plot of avalanche breakdown voltage vs critical field for 4H-SiC DIMOSFET for different values of $N_o$ . . . . .	32
4.9	Plot of punch through breakdown voltage vs device height(region I) for different values of $N_o$ . . . . .	32
4.10	Plot of punch through breakdown voltage vs avalanche breakdown voltage for different values of $N_o$ . . . . .	33
4.11	Plot of on-state depletion width vs specific-on resistance voltage for different values of $N_o$ . . . . .	33

# List of Tables

1.1	Comparison between Silicon and Silicon Carbide poly types . . . . .	3
4.1	Results of currents, voltages, $R_{on-sp}$ and $P_D$ for doping $10^{12}$ - $10^{15}$ in region I and $10^{15}$ - $10^{16}$ in region II, $h_1=0.0272\text{cm}$ , $N_{eff}=2.85 \times 10^{17}$ , $\alpha_1=3.674 \times 10^{16} \text{ cm}^{-4}$ . .	29
4.2	Results of currents, voltages, $R_{on-sp}$ and $P_D$ for doping $2 \times 10^{12}$ - $10^{15}$ in region I and $10^{15}$ - $10^{16}$ in region II, $h_1=0.0258\text{cm}$ , $N_{eff}=2.65 \times 10^{17}$ , $\alpha_1=3.87 \times 10^{16} \text{ cm}^{-4}$	29
4.3	Results of currents, voltages, $R_{on-sp}$ and $P_D$ for doping $2.5 \times 10^{12}$ - $10^{15}$ in region I and $10^{15}$ - $10^{16}$ in region II, $h_1=0.0253\text{cm}$ , $N_{eff}=2.58 \times 10^{17}$ , $\alpha_1=3.94 \times 10^{16} \text{ cm}^{-4}$	29
4.4	Results of currents, voltages, $R_{on-sp}$ and $P_D$ for doping $10^{13}$ - $10^{15}$ in region I and $10^{15}$ - $10^{16}$ in region II, $h_1=0.0223\text{cm}$ , $N_{eff}=2.15 \times 10^{17}$ , $\alpha_1=4.44 \times 10^{16} \text{ cm}^{-4}$ . .	30
4.5	Results of breakdown voltages, $V_{BAV}$ for a punch through breakdown, $V_{BPT}$ set to 10kV . . . . .	34
4.6	Results of currents, voltages, $R_{on-sp}$ and $P_D$ for doping $10^{15}$ , $h=0.0073\text{cm}$ for uniformly doped drift region . . . . .	35
4.7	Results of currents, voltages, $R_{on-sp1}$ and $P_D$ for doping $10^{15}$ - $10^{19}$ in region I and $10^{19}$ - $10^{20}$ in region II, $h_1=0.000342\text{cm}$ , $N_{eff}=1.09 \times 10^{18}$ , $\alpha_1=1.54 \times 10^{21} \text{ cm}^{-4}$ for distorted profile . . . . .	35
4.8	Results of calculation of breakdown voltages, $V_{BAV}$ for a set punch through voltage, $V_{BPT} = 5\text{kV}$ for uniformly doped profile and distorted Gaussian profile	35

# List of Symbols

$E_c$	Critical electric field
$W_j$	Depth of p-body
$I_{DS}$	Drain current
$N_{eff}$	Effective carrier concentration
$\mu_n$	Electron mobility
$n_i$	Intrinsic carrier concentration
$L$	Length of channel
$\tau_e(\tau_p)$	Lifetime of an electron(lifetime of a hole)
$W_d$	On-state depletion region width
$C_{ox}$	Oxide capacitance
$v_{sat}$	Saturated drift velocity of electron
$m$	Straggle value of Ion-Implantation
$W_t$	Total device height
$W$	Width of channel

# List of Abbreviations

ACCUFET	Accumulation mode Field Effect Transistor
AlN	Aluminium Nitride
DIMOSFET	Double Implant Metal Oxide Semiconductor Field Effect Transistor
FET	Field Effect Transistor
GaN	Gallium Nitride
JFET	Junction Field Effect Transistor
MBE	Molecular Beam Epitaxy
MOS	Metal Oxide Semiconductor
MOSFET	Metal Oxide Semiconductor Field Effect Transistor
$P_D$	Power dissipation
$R_{on-sp}$	Specific on-resistance
Si	Silicon
SiC	Silicon Carbide
UMOSFET	U - shape MOSFET
$V_{BAV}$	Avalanche breakdown voltage
$V_{BPT}$	Punch through breakdown voltage
$V_T$	Threshold voltage
WBG	Wide Band Gap
eV	electron Volt

# Contents

<b>ACKNOWLEDGEMENT</b>	<b>i</b>
<b>ABSTRACT</b>	<b>ii</b>
<b>LIST OF FIGURES</b>	<b>iii</b>
<b>LIST OF TABLES</b>	<b>iv</b>
<b>LIST OF SYMBOLS</b>	<b>v</b>
<b>LIST OF ABBREVIATIONS</b>	<b>vi</b>
<b>1 INTRODUCTION</b>	<b>1</b>
1.1 Preamble . . . . .	1
1.2 Polytypes of Silicon Carbide . . . . .	2
1.3 Power MOSFETs . . . . .	4
1.4 Device structure and analysis of 4H-SiC DIMOSFET . . . . .	5
1.4.1 Forward Conduction Characteristics . . . . .	6
1.5 Motivation . . . . .	10
1.6 Scope and Objectives . . . . .	10
1.7 Organization of the Report . . . . .	11
<b>2 LITERATURE REVIEW</b>	<b>13</b>
<b>3 A COMPARATIVE STUDY ON UNIFORMLY DOPED AND LINEARLY GRADED PROFILE</b>	<b>18</b>
3.1 Introduction . . . . .	18
3.2 Vertical DIMOSFET with Uniformly Doped Profile . . . . .	18
3.3 Vertical DIMOSFET with Linearly Graded Doping Profile . . . . .	19
3.4 Discussion of results . . . . .	20

<b>4 ANALYSIS AND DESIGN OF 4H-DIMOSFET USING DISTORTED GAUSSIAN DOPING PROFILE</b>	<b>21</b>
4.1 Introduction . . . . .	21
4.2 Vertical DIMOSFET with Distorted Gaussian Profile in the Drift Region . . . . .	22
4.3 Theoretical Analysis . . . . .	22
4.3.1 Calculation of effective carrier concentration( $N_{eff}$ ) . . . . .	24
4.3.2 Calculation of Channel Voltage ( $V_{ch}$ ) . . . . .	26
4.3.3 Critical Field, $E_c$ and Breakdown Voltage Calculations . . . . .	27
4.4 Calculations and Related Graphs . . . . .	28
4.4.1 Calculation of Breakdown Voltages( $V_{BAV}$ and $V_{BPT}$ ) . . . . .	34
4.5 Comparison between uniformly doping profile and distorted Gaussian profile . . . . .	34
4.6 Discussion of Results . . . . .	36
<b>5 CONCLUSION AND FUTURE SCOPE OF WORK</b>	<b>37</b>
<b>REFERENCES</b>	<b>40</b>
<b>A APPENDIX</b>	<b>44</b>

# Chapter 1

## INTRODUCTION

### 1.1 Preamble

Silicon carbide (SiC) has proved to be a potential game changer in power semiconductor area. Properties such as large breakdown electric field(critical field), high saturated drift velocity, small dielectric constant, wider band gap and high thermal conductivity makes SiC an attractive alternative for fabricating power semiconductor devices which still is primarily dominated by Silicon. However, research works pertaining to SiC devices has not received much attention due to lack of reproducible techniques to grow single crystals and epilayers of the same. Silicon Carbide (SiC), as discussed above has superior properties(both electrical and mechanical) for power devices compared to Silicon. In recent times there is a need for power devices to operate at higher power and temperature levels which has pushed for more research in the area of this material.

Among the electrical properties, wide band gap is an important one while modeling a power device; a wider band gap enables the designer to design highly dense devices with high voltage withstanding ability. Silicon Carbide belongs to the class of wide band gap semiconductors, while semiconductors like Si, Ge and GaAs belong to narrow band gap semiconductors. Among the Wide Band Gap semiconductors, Gallium Nitride(GaN) and Aluminium Nitride(AlN) have higher band gap than SiC(3.4eV and 6.2eV respectively against 3.2eV of 6H-SiC), however the primary reason why SiC is more preferred over former is because of it's ability to form high quality native oxide(i.e, SiO<sub>2</sub>). The large Si-C bonding energy also makes SiC resistant to chemical attack and radiation. High thermal conductivity is another important feature of SiC material. Thermal conductivity is a measure of the rate at which heat passes through a specified material. The higher thermal conductivity( $\approx 4.9$  W/m-K for 4H and 6H-SiC) ensures less packaging and heat sinking cost[29],[30]. These are the few among various reasons why

SiC is pushed over Si for power devices.

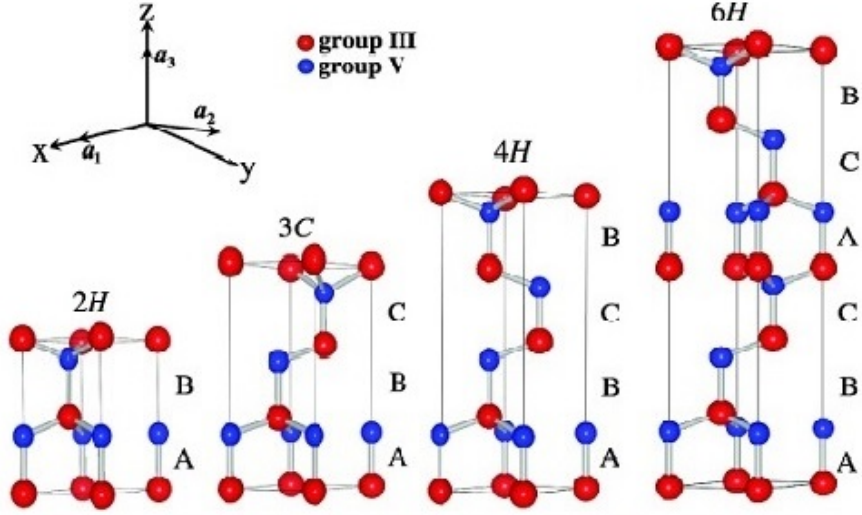
## 1.2 Polytypes of Silicon Carbide

Polytypes refer to the different crystal/lattice structures of the same element. For Silicon Carbide there are nearly 200 polytypes which differ from each other only in their stacking sequences (stacking sequence refers to how each layer of Silicon and Carbon are aligned/stacked over each other.).

In Figure-1.1, the stacking sequence is shown for the three commonly used polytypes of SiC, viz. 3C, 4H and 6H-SiC (devices based on only these three polytypes have been successfully developed so far). In 4H, if a layer is designated as A, then next layer is not centered above A due to tetrahedral configuration of Si, this layer is named B. The third layer is centered above A but is complemented, this layer is  $\bar{A}$  and the fourth layer is not centered above A or B which appears as complemented thus, assigned C. Hence, the stacking sequence in 4H-SiC is  $AB\bar{A}\bar{C}\dots$  as shown in Figure-1.1. On the similar lines, 6H-SiC has the stacking sequence as  $ABCBA\bar{A}\bar{C}\dots$  [30]. The change in the lattice structure affects the electrical properties a lot. Some of these differences are listed in Table-1.1.

Among those many polytypes of Silicon Carbide, 3C, 4H and 6H-SiC electronic devices presently are the most widely experimented as they are readily available and also due to high quality single crystal wafers unlike the remaining polytypes. These reasons have helped 3C, 6H and 4H-SiC polytypes to be manufactured in bulk wafer form which has helped SiC to emerge as one of the relatively mature wide-band semiconductor technologies [27],[30]. As per the discussions above, it has been proved that Silicon Carbide is indeed a better alternative for Silicon in power device. However, there are many crucial issues like single crystal growth, device fabrication issues etc. that have to be looked upon before SiC can be accepted as a reliable material for semiconductor power industry [25]. Table 1.1 lists some electrical properties of the SiC polytypes and compares them with that of Si.

As seen in Table 1.1, the band gap of SiC is almost thrice of Silicon, mobility is comparable to that of Silicon (4H and 3C polytypes) while thermal conductivity is way higher than Silicon. Intrinsic concentration of SiC is almost negligible compared to Si. This lower value is due to higher band gap and this results in lower leakage as temperature rises, making it more favored material than silicon for high power applications.



**Figure 1.1:** Stacking sequence of the most common polytypes of SiC[30]

**Table 1.1:** Comparison of electronic properties among Si and different polytypes of SiC[30]

Parameters	Si	6H-SiC	4H-SiC	3C-SiC
Bandgap (eV)	1.1	3	3.26	2.2
Breakdown field @ $10^{17}\text{cm}^{-3}$ (MV/cm)	0.6	3.2	3	1.5
Electron mobility @ $10^{16}\text{cm}^{-3}$ ( $\text{cm}^2/\text{V}\cdot\text{s}$ )	1100	370	800	750
Saturated electron drift velocity (cm/s)	$10^7$	$2 \times 10^7$	$2 \times 10^7$	$2 \times 10^7$
Intrinsic concentration, $n_i$ ( $\text{cm}^{-3}$ )	$1.5 \times 10^{10}$	$2.3 \times 10^{-6}$	$8.2 \times 10^{-9}$	$6.92 \times 10^{-9}$
Thermal conductivity (W/cm-K)	1.5	4.9	4.9	5

Among the polytypes of Silicon Carbide, 4H and 6H-SiC electronic devices presently are the most promising because of the availability and quality of reproducible single-crystal wafers. The availability of 6H-SiC and 4H-SiC polytypes in bulk wafer form has helped SiC to emerge as one of the relatively mature wide-band semiconductor technologies[31]. SiC is a material with immense potential for use in hetero structure electronic devices, which take advantage of differing band gaps, carrier mobilities etc. However, there are many crucial issues, crystal growth and device fabrication to name a few, that have to be addressed before SiC-based devices and circuits can be scaled up and reliably incorporated into electronic systems. The most important issue being controlled and repeatable doping in SiC device structures.

### **1.3 Power MOSFETs**

Compared to their bipolar counterpart, MOSFETs are undoubtedly superior with faster switching time, high input impedance(which leads to a simpler driver circuitry), lower leakage devices(as they are majority carrier device)[31]. MOSFETs are voltage-controlled device while a bipolar transistor is a current controlled device. The presence of majority carriers makes MOSFET a faster switching device than bipolar devices(BJT), which is useful during hard switching. However, carrier mobility decreases with increase in temperature(due to impurity vibrations at higher temperature) which results in slower device at elevated temperatures. This makes MOSFET more resistive at higher temperatures but they are immune to thermal-runaway problem as experienced by bipolar devices. This property eliminates the need for additional heat sink circuitry requirements[27][30].

There is no doubt that power switches are the heart of all power electronic systems. Many reasons have made power electronic systems affordable in large number of applications; the increased power capabilities, ease of control and reduced costs of power switches to name a few. Thyristors and bipolar transistors were the first power switches. Until the late 1970s, the rating of these devices grew steadily. It was around this time that the first power MOSFET was introduced. Since then, power MOSFETs based on silicon have drastically improved and become the dominant device technology since 1980s for many applications(and is continuing so till now). The reasons why MOSFETs are preferred over their bipolar counterparts are[22][30]:

- MOSFETs have relatively simpler gate drive circuitry due to high input impedance. As the input impedance is high, just simple integrated circuits can be used to control the drive of gate as it requires very low gate current.
- MOSFETs possess higher switching speeds than bipolars as they are majority carrier

device. This absence of minority charge carriers also lead to low leakage at higher temperature for power devices.

- MOSFETS have better ruggedness and safe operating forward biased area, which eliminates the need of protective snubber circuits during hard switching, unlike bipolar transistors.
- As the resistivity of majority carriers increases with temperature, the thermal runaway behavior is avoided in MOSFETs, which eliminates the need for additional heat sinking circuitry.

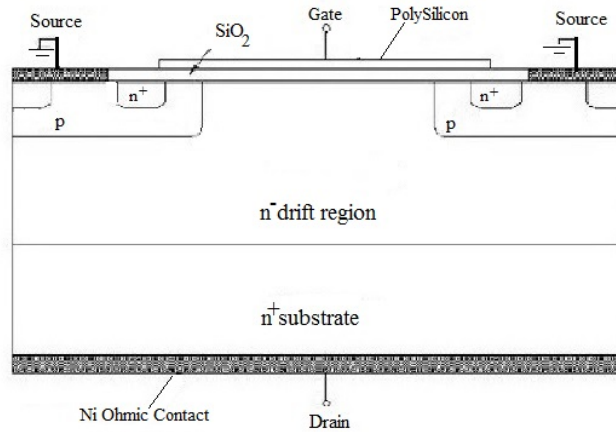
## 1.4 Device structure and analysis of 4H-SiC DIMOSFET

A power MOSFET is a voltage driven device whose gate terminal is electrically isolated from its Silicon Carbide body by a thin layer of silicon dioxide ( $\text{SiO}_2$ ). Due to very few minority charge carriers the speed of operation of MOSFET is higher. DIMOS(Double Implant MOS) transistors are common in silicon power device technology where the p-base and  $n^+$  source regions are formed by diffusion of impurities through a common mask opening. However, owing to very low diffusion constant for SiC, impurity diffusion is impractical, the favored doping technique is ion implantation.

The first attempt in doping SiC using ion implantation was done by Purdue group. The implantation required two separate masks for doping p-base and  $n^+$  regions respectively, as shown in Figure-1.2, this is the reason for the nomenclature Double Implant. The device fabrication is started with an n-type epitaxial layer grown on a heavily doped  $n^+$  substrate. The channel is formed by the difference in lateral extension of the p-base and  $n^+$  source regions produced by their diffusion cycles. Both regions are self-aligned to the left-hand-side and right hand- side of the gate region during ion-implantation to introduce the respective dopants.

Without the application of a gate bias, a high voltage can be supported in the DIMOS structure when a positive bias is applied to the drain. In this case, a reverse bias junction is formed between the p-base and the n-drift region. Drain current that flows in the DIMOSFET structure is induced by the application of a positive bias to the gate electrode. This produces an inversion layer at the surface of the p-base region under the gate electrode. This inversion layer channel provides a path for transport of electrons from the source to the drain. After transport from the source region through the channel, the electrons enter the n-drift region at the upper surface of the device structure. They are then transported through a relatively narrow JFET region located

between the adjacent p-base regions within the DIMOSFET structure. After being transported through the JFET region, the electrons enter the n-drift region. The current spreads from the relatively narrow JFET region to the entire width of the cell cross-section. This non-uniform current distribution within the drift region enhances its resistance making the internal resistance of the DIMOSFET structure larger than the ideal specific on-resistance of the drift region[23].

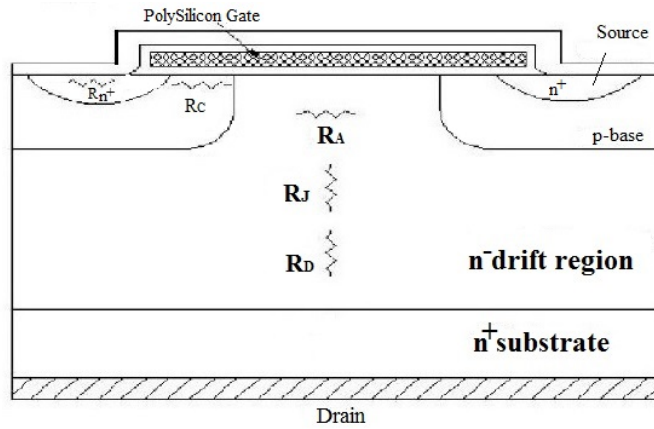


**Figure 1.2:** Structure of DIMOSFET[29],[30]

### 1.4.1 Forward Conduction Characteristics

The current flow in DIMOSFET can be understood by analyzing the device structure. Unlike planar structure, DIMOSFETs are vertical structures with source terminal located at the top and drain terminal at the bottom. The vertical structure also enhances the current flowing along the device. The channel formation (between the n<sup>+</sup> source and the p-body) is due to the potential applied at the gate terminal. The channel is formed in the gap between p-base and n<sup>+</sup> diffusion regions. The current flow in the power DIMOSFET during forward conduction is limited by the total resistance between the source and drain. The resistance is limited and comprise of many components. These several components are as shown in Fig.1.3. Here,  $R_{n^+}$  is the contribution from n<sup>+</sup> source diffusion region,  $R_C$  is the channel resistance,  $R_A$  is the accumulation layer resistance,  $R_D$  is the drift region resistance,  $R_S$  is the substrate resistance and the portion of the drift region that comes to the upper surface between cells that contribute  $R_J$  that is enhanced at a higher drain voltage due to pinch-off action of depletion layer extending from the p-base regions due JFET action.

The on-resistance of the device is the total resistance between the source and drain terminals in the on-state [23],[30],[32]. This parameter is responsible for the current rating of the device

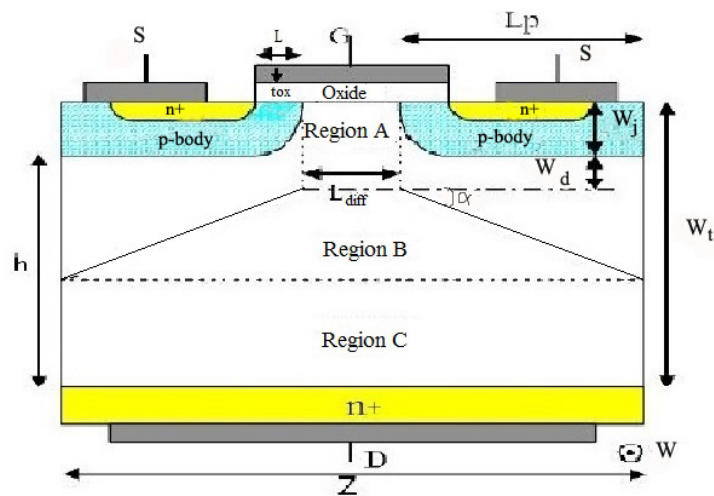


**Figure 1.3:** On-state resistance distribution in DIMOSFET[29],[33].

and consequently the power dissipation. The cell structure with each component of the specific on-resistance ( $R_{on-sp}$ ) is as shown in Figure-1.3. The application of positive gate bias results in a current flow between drain and source through the n-drift region and conductive channel. The conductivity of the channel is modulated by the gate bias voltage and the current flow is determined by the resistance of various resistive components as shown in Figure-1.3. The total specific on-resistance ( $R_{on-sp}$ ) is determined as[29],[30]:

$$R_{on-sp} = R_{n+} + R_C + R_A + R_D + R_J \dots (1.1)$$

Under lower breakdown voltages, all the resistances come into picture while at higher breakdowns the drift region resistance ( $R_D$ ) dominates.



**Figure 1.4:** Cross section of DIMOSFET[30].

In a power MOSFET, the blocking voltage is supported across the drift layer, and thus the drift region resistance is considered to be the minimum possible theoretical limit for the on-resistance of a MOSFET. For an ideal DIMOSFET, the resistances associated with the n<sup>+</sup> source, the n-channel, the accumulation region and the n<sup>+</sup> substrate are assumed to be negligible and the specific on-resistance of the power MOSFET is determined by the drift region only. This assumption is not accurate at lower breakdown voltages where the drift region resistance  $R_D$  is comparable to the other resistive components and these resistances should be included in calculating  $R_{on-sp}$ . However, at higher breakdown voltages,  $R_D$  is significantly higher than other resistances and  $R_{on-sp}$  can be approximated to  $R_D$  [30].

Fig.1.4 depicts the direction of current flow in the power MOSFET. The mathematical formulas pertaining to various resistance values is given as follows. If the linear cell is considered with 1-cm extension perpendicular to the cross section as shown in Figure-1.3, the resistance per square centimeter due to the n<sup>+</sup> region is given by[30]:

$$R_{n+} = 1/2\rho_{ON}L_E(L_G + 2m) \dots (1.2)$$

where,  $\rho_{ON}$  is the sheet resistance of n+ diffusion and parameters  $L_E$  and  $L_G$  are the emitter and gate lengths respectively,  $2m$  is the cell diffusion window. The channel resistance per square centimeter for the linear cell is given by [25],[30]:

$$R_C = \frac{1/2L(L_G + 2m)}{(\mu_n C_{ox}(V_G - V_T))} \dots (1.3)$$

The resistance of the accumulation layer  $R_A$  determines the current spreading from the channel into drift region. The accumulation layer resistance is dependent on the charge in the accumulation layer and the mobility for free carriers at the accumulated surface. For the linear cell geometry, the accumulation layer resistance per square centimeter is [25],[30]:

$$R_A = \frac{K(L_G + 2m)(L_G - 2x_p)}{(\mu_n C_{ox}(V_G - V_T))} \dots (1.4)$$

where,  $K$  is a factor introduced to account for the two dimensional nature of the current flow from the accumulation layer into the bulk. The resistance of drift region between the p-base diffusion can be calculated if the voltage drop along the vertical direction is neglected. Under the assumption that the depletion layer width of the p-base-n-drift layer junction is negligibly small and the current is flowing uniformly down from the accumulation layer into the JFET region. The resistance of the JFET region can be analyzed as a resistance with increasing cross

section when proceeding downward from the surface and is calculated by[25],[30]:

$$R_J = 2\rho_D(L_G + 2m)\left[\frac{1}{\sqrt{1 - (2x_p/L_G)^2}} \tan^{-1}(0.414) \sqrt{\frac{(L_G + 2x_p)}{(L_G - 2x_p)} - \frac{\pi}{8}}\right] \dots (1.5)$$

where,  $\rho_D$  is the resistivity of the drift region. For the current spreading at an angle as shown in Figure-1.3, the drift region resistance per square centimeter is given by [30]:

$$R_D = \rho_D \frac{(L_G + 2m)}{\tan \alpha} \ln[1 + 2(h/a) \tan \alpha] \dots (1.6)$$

where, dimensions h and  $\alpha$  are indicated in Figure-1.3.

A good approximation for  $\alpha_1$  in theory is provided by[29],[30]:

$$\alpha_1 = 28^\circ - (h/a) \text{ if } h \geq a \dots (1.7)$$

$$\alpha_1 = 28^\circ - (a/h) \text{ if } h \leq a \dots (1.8)$$

Finally, the equation for power dissipation for 50% duty cycle can be evaluated from the equation given as[18]:

$$P_D = \frac{1}{2}((J_{ON})^2 AR_{on} + J_L AV_B) \dots (1.9)$$

where,  $J_{ON}$  is the on-state current density,  $J_L$  is the leakage current density,  $V_B$  is the breakdown voltage and A is the cross-sectional area of the device. The leakage current density ( $J_L$ ) is given as[18]:

$$J_L = \frac{en_i W}{\tau_e} + e\left(\frac{D_h}{\tau_h}\right)^{\frac{1}{2}} \frac{n^2}{N_d} \dots (1.10)$$

where,  $n_i$  is the intrinsic carrier concentration ( $\text{cm}^{-3}$ ), W is the width of the space-charge region (cm),  $\tau_e$  is the lifetime of electrons in the space-charge region (s),  $N_D$  is the concentration of donor atoms ( $\text{cm}^{-3}$ ),  $D_h$  is the diffusion constant of holes in the n-type region ( $\text{cm}^2/\text{s}$ ) and  $\tau_h$  is the lifetime of holes(s). Thus, the equations from eq.(1.1) to eq.(1.10) describes all the equations required for theoretical analysis.

## 1.5 Motivation

Silicon MOSFETs are already the dominant technology in simpler low power, high switching electronic devices. However, the relative lower breakdown field strength of Silicon limits their usage in applications required to operate at higher temperatures (500° and above). The advantages of SiC material properties, in particular breakdown field, makes SiC MOSFETs a very promising candidate for fast switching devices required at higher temperature. SiC has much lesser power dissipation at higher temperatures than Silicon due to less specific on-resistance (which is almost 100-200 times less). Lower thermal minority carrier generation implies lower leakage currents and device operation at higher temperatures, arising from self heating due to power dissipation is more tolerable. Moreover, the thermal conductivity of SiC is three times higher than Si which eliminates the need for additional circuits for heat sinking purpose.

## 1.6 Scope and Objectives

### Scope:

Due to excellent physical and electrical properties such as high breakdown electric field, wide band gap, high thermal conductivity and high electron saturation velocity, silicon carbide offers great potential for development of high temperature, high power and high voltage devices. Significant progress in SiC power MOSFETs have been demonstrated with the fabrication of UMOS, DIMOSFET (Double Implant MOSFET), triple implanted vertical MOSFET and accumulation mode MOSFET (ACCUFETs)

### Objective:

The objective of this thesis work is to analyze one of the poly types of SiC (4H-SiC) for higher breakdown voltage and lower specification on-resistance (and consequently lower power dissipation). The profile considered for the drift region is a distorted Gaussian profile. The objectives are summarized as:

- Study and Analysis of Double Implanted MOSFET on 4H-SiC wafer.
- Analysis and Design of Double Implanted MOSFET with optimum power dissipation for large breakdown voltages.

- Analysis and Design of 4H-DIMOSFET with distorted Gaussian doping profile of the drift region .

## 1.7 Organization of the Report

In this report, the **Chapter 1** is an introduction to Silicon Carbide (SiC) material and the working of the device(DIMOSFET). Recent advancements in this field is discussed.A brief discussion on different polytypes of SiC and and on why only 4H, 6H and 3C polytypes are considered in power devices is included.SiC is superior to Silicon(Si) due to properties like higher critical electric field, wider band gap, higher thermal conductivity to name few. A comparison table is drawn to indicate the differences among different polytypes of SiC and comparing it against Si. The device under consideration is DIMOSFET,which unlike the planar MOSFET devices is a vertical device. Vertical devices have the source and drain terminals located at the top and bottom respectively. This arrangement ensure higher current flow along the device height.

**Chapter 2** covers the reviews of literature that were used as reference during the research work in chronological order. The first MOSFET came up in 1970s and power MOSFETs based on silicon has been the dominant technology since 1980s. Many of the works carried out on SiC power devices have been implemented using uniformly doped or linearly graded profile. There has been almost no reference to any work wherein non-linear doping profile has been used. All the limitations and advancements over the years has been mentioned in detail in this chapter.

**Chapter 3** covers a comparative study between the two widely used doping profiles in power industries for power MOSFETs namely, uniformly doped profile and linearly graded profile. Results from previous work based on uniformly doped profile and linear one have been compared against each other. A linearly graded profile gives smaller device height than a uniform doped one for the same breakdown voltage. There is a limit to the breakdown voltage that can be achieved using uniform doping profile and there exists a trade-off between power dissipation and breakdown voltage in such profiles. A non-linear profile, like Gaussian and complementary error, overcomes this trade-off by the virtue of their non-linearity.

This is followed by **Chapter 4** which carries the analysis of the distorted Gaussian doping profile. The distorted Gaussian profile (a more practical non-linear doping profile) is approximated by a dual-slope linearly graded profile. The analysis includes calculation of effective carrier concentration in the drift region( $N_{eff}$ ), channel voltages. The critical field and the break-

down voltages, avalanche breakdown and punch through breakdown voltage, is then derived for the above profile in this chapter. This is followed by the required calculations and the respective plots. The various plots includes current density versus power dissipation, current density versus on-state depletion width for different peak carrier concentration, current density versus drain-source voltage, effective carrier concentration versus concentration gradient, avalanche breakdown voltage versus critical electric field, avalanche breakdown voltage versus device height, avalanche breakdown voltage versus punch through breakdown voltage and on-state depletion width versus specific on-resistance. The chapter is concluded by a comparison between uniformly doping profile and distorted Gaussian profile. Finally, all the discussion of the results based on plots and tables in the work is discussed in **Chapter 5**. This chapter also contains the future scope of the work. All the references used in the work are enlisted in the final chapter. The Appendix section carries the detailed derivation for the effective carrier concentration for the two-step linearly graded doping profile for the device structure discussed in the report.

# Chapter 2

## LITERATURE REVIEW

The rise of Silicon Carbide was in the late 1980s, when it was observed that power devices based on Silicon were reaching its theoretical limits and that limit could be extended significantly by using material with higher breakdown field[3]. However, it was only in 1994 that the MOS characteristics of silicon carbide was explained in a paper by Brown et al[5]. This paper also came up with an explanation characterizing the SiC/SiO<sub>2</sub> interface states which was due to rejection of n-type dopants and incorporation of p-type dopants into oxide during growth. However, both Si and SiC showed the same impurity redistribution during the thermal oxidation. It was in late 1980s that the first SiC based MOSFET was designed but the first SiC based power MOSFET came up only in 1994 which were basically vertically trenching MOSFET or UMOSFET[6]. But UMOSFETs had two major issues[7-8]:

- Due to U shape of the trenches, at the corners fringing of field happened leading to failure of gate oxide at higher drain voltages thereby limiting the breakdown voltage much lower than the theoretical value.
- Lower specific on-resistance (hence higher power dissipation) due to low inversion layer mobility of carriers along the trench sidewalls.

In 1995, a SiC U-MOSFET with maximum breakdown voltage of only 260V was designed. B.Jayant Baliga in his journal[9] described the applications of different SiC devices that were commercially available in the year 1996. More and more testing and fabrication of many SiC MOSFETs started in the years that followed. Year 1998 saw the mention of 4H-SiC UMOSFETs and DIMOSFETs with measured blocking voltages of 1400 V and 900 V respectively in [10]. The paper also highlighted carrier mobility, high interface state density, the difficulties of forming high quality oxide on the sidewalls of the vertical trenches(for UMOSFETs). An year later(in 1999), the numerical and theoretical analysis of 6.5kV SiC based JFET and MOSFET

came up[11]. The same year elucidated the usefulness of Silicon Carbide in device applications. For the above said device, the n-doped drift region was  $60\mu\text{m}$  long with  $2 \times 10^{15}\text{cc}$  doping concentration and the gate and source diffusion were  $0.6\mu\text{m}$  apart. The gate oxide thickness was  $0.2\mu\text{m}$  for the simulated structure. All the simulations and optimizations of the above device structure were done in MEDICI simulator.

The different MOSFET parameters like density of interface states, intrinsic mobility reduction factor, inversion layer mobility etc. were studied and examined in depth for 6H SiC enhancement mode n channel MOSFET[12]. It was found that with increase in substrate doping concentration, there was a reduction in inversion layer mobility. There was an analogy to silicon devices with regard to dependencies of inversion layer on substrate doping. However, these dependencies are modified by high concentration of interface states near the edge of conduction band. The concept of interface states was described in detail in[13] including the effect of these states on the carrier mobility for 4H and 6H-SiC. This study was necessary to understand why carrier mobility decreases for 4H and 6H-SiC and not for 3C-SiC. The same year J.Wang et al in his paper[14] came up with a trade-off between breakdown voltage and drift region on-resistance and evaluated the high blocking capacity of 4H-SiC DIMOSFET while compromising for low on resistance.

The year 2000 saw the characterization of Silicon Carbide epitaxial channel MOSFETs. These MOSFETs were fabricated on 6H SiC substrates with n+ epitaxial source and drain terminals. Under 50% channel donor impurity concentration, a buried channel mobility of  $230\text{ cm}^2/\text{V-s}$  and accumulation channel mobility of  $45\text{ cm}^2/\text{V-s}$  were extracted[16]. The future of SiC switching devices was discussed in[18]. In the same year, for the first time,  $\text{N}_2\text{O}$  grown oxide on both n and p-type 6H SiC wafers was successfully demonstrated and fabricated. This  $\text{N}_2\text{O}$  grown oxide showed better oxide qualities and improved SiC/ $\text{SiO}_2$  interface characteristics. This  $\text{N}_2\text{O}$  oxidation technique showed highly reliable and improved quality SiC MOSFETs. In the following year 4H-SiC RF MOSFET was fabricated for the first time. The enhanced performance of this device was due to two-metal layer process which optimized the complementing requirements of acceptable value of inversion- layer mobility and low contact resistance. The enhanced performance included a breakdown voltage of 950V attained in MOSFET with specific on-resistance of  $24\Omega\text{-mm}^2$ . Electron mobility model for three of the main SiC poly types, namely 3C, 4H and 6H were developed.

A. Mihaila et al in the paper[17] presented a systematic analysis of breakdown mechanisms

in silicon carbide MOSFET and JFET. Trench technology was used for the MOSFET and the devices designed for 1.2kV were simulated and optimized using MEDICI and ISE TCAD software packages. However many drawbacks in SiC trench MOSFET (like gate oxide breakdown, low channel mobility and the tight trade-off between the punch-through premature breakdown and the threshold voltage in the channel) were eliminated by using the SiC JFET. In the years that followed many developments took place in SiC MOSFETs. Specific behaviour of interface states were measured which indicated that at 0.1eV below the conduction band the interface state density decreased from  $2 \times 10^{13}$  to  $2 \times 10^{12}$  eV<sup>-1</sup> cm<sup>-2</sup> following annealing in N<sub>2</sub>O for 2 hours[19]. A 2D drift-diffusion based simulator was designed for SiC in the same paper. Even by this time fabricating SiC devices had their limitations which included material quality, ion implantation, the SiC-SiO<sub>2</sub> interface and the thermal stability of contacting systems which called for further work[20]. Evaluations were done for SiC based Schottky and PIN diodes against their Silicon counterparts.

In the year 2002, a breakdown of only 760V have been attained from 6H- DIMOSFET devices [21]. Pyrogenic re-oxidation, could reduce the specific on-resistance and thereby significantly reduce the power dissipation along the device. In the same year, an analytical model of a SiC MOSFET came up in[23]. In this paper, a semi-empirical relation for carrier mobility dependency on dopant concentration, electric field intensity and temperature was formulated which was based on experimental results. Based on this dependencies, apt analytical-physical model of V-I characteristics, trans-conductance and conductance were developed. These proposed models led to the design of a simulation algorithm which then was used to simulate the MOSFET performance. In the same year, the first SiC MOSFET with breakdown in kV range was demonstrated - a 10A, 2.4kV 4H SiC DIMOSFET. This MOSFET, at room temperature, showed a peak channel mobility of 22 cm<sup>2</sup>/V-s, a threshold voltage of 8.5V and exhibited a specific on-resistance of 42 mΩ-cm<sup>2</sup>. At a temperature of 200degC increased to a value of 85 mΩ-cm<sup>2</sup>. However, a stable value for avalanche breakdown of 2.4kV was observed. The MOS channel length, in the paper, as defined by the p-well and n-implant was 1.5μm. In this vertical device, the charge carriers flow laterally from the n-source through the channel on the implanted p-well and then vertically through the drift region which is spread at an angle. The device dimensions chosen for a breakdown of 2kV was a 20μm thick drift region with a doping concentration of  $2.5 \times 10^{13}$  cc. Apart from all the above specifications, the paper also pointed out that these devices were capable of high breakdown, low loss switching and high frequency applications.

One of the main reasons why SiC based devices have found their place in power industry

is because of their higher breakdown voltage than Si. To extend its viability in high power applications, its required to enhance the breakdown voltage. Numerical device simulations on 4H-SiC vertical MOSFETs were presented in [24-25]. Reverse blocking voltage, onresistance and threshold voltage were the main focus in these simulations. The simulated model had a gate oxide thickness of 50nm, p-well depth as well as channel length of micron and a source depth of  $0.2\mu\text{m}$ . As shown in [26], breakdown of 910V could be attained from one-step field plate termination and 1100V from the embedded mesa with step field plating. Reliability and performance became the next concern for SiC devices. These properties were then the rising demand in semiconductor material, since with the rising breakdown voltages, higher operating frequencies and higher temperatures the chances of the device wearing out and breaking down would increase.

In the paper[27], the same were studied for advanced power applications. All these years, the analysis and models were developed for uniformly and linearly doped profiles only. Since these profiles could be easily developed and analyzed. However, there was a major limitation with these profiles; either power dissipation or breakdown voltage could be optimized at a time. To overcome this limitation, non-linear doping profiles were needed to be considered. First theoretical analysis to consider non-linear doping profiles on 6H-SiC DIMOSFET (like Gaussian, complementary error profile) appeared in [29-30]in the year 2008. The advantage these non-linear profiles seemed to have over uniformly doped and linearly graded profile was that it could give higher breakdown while keeping lower values for parasitic resistances. The profile had a very low value of impurity concentration near the p-well and n-drift region while a higher concentration at the lower end of the device. However, practical implementation of such a structure is yet to be seen. In paper[30], the power dissipation for the linearly graded profile evaluated at a current density of  $1000\text{ A/cm}^2$  was found to be minimum, whereas a maximum breakdown voltage of 20kV was found for the Complementary Error Function profile. Device structure of a 6H-SiC vertical Double-Implanted MOSFET (DIMOSFET) to provide a high breakdown voltage of about 10 kV and a low power dissipation for a rise in device temperature of  $600\text{ }^\circ\text{C}$  was studied in[29]. For the same temperature, the device showed optimum doping levels of the drift region lying between  $5\times 10^{13}\text{ cm}^{-3}$  and  $5\times 10^{15}\text{ cm}^{-3}$  for a breakdown voltage of 10 kV. Analysis results also showed that the devices with Complementary Error Function profile had smaller device height for the same breakdown voltage as compared to linearly graded or uniformly doped profiles.

The recent years saw a major leap in the SiC device technology with a recently(year 2014) developed 2nd generation, large-area ( $56\text{ mm}^2$  with an active conducting area of  $40\text{ mm}^2$ ) 4H-

SiC DMOSFET, which reliably blocked 1600 V with very low leakage current under a gate-bias ( $V_G$ ) of 0 V at temperatures up to 200. The device also exhibited a low on-resistance ( $R_{on-sp}$ ) of 12.4 m $\Omega$  at 150 A and  $V_G$  of 20 V. DC and dynamic switching characteristics of the SiC DMOSFET were also compared with a commercially available 1200 V/ 200A rated Si trench gate IGBT[33]. This paper also presented a comprehensive study on intrinsic reliability of this 2nd generation SiC MOSFET to build consumer confidence and to achieve broad market adoption of power switch technology. In another paper[34] to obtain widespread usage in SiC DMOSFETs, their long-term operational ability to handle the stressful transient current and high temperatures common in power electronics were further verified. The long-term reliability of a single 4H-SiC DMOSFET, the effects of extreme high current density were evaluated. The 4H-SiC DMOSFET, which was studied, had an active conducting area of 40 mm<sup>2</sup>, and ratings of 1200 V and 150 A. The device was electrically stressed by hard switching transient currents in excess of four times the given rating ( $\geq 600A$ ) corresponding to a current density of 1500 A/cm<sup>2</sup>. Periodically throughout testing, several device characteristics including  $R_{on-sp}$  and  $V_{GS(th)}$  were measured. The hard switching was done 500000 times, the first 16000 switching cycles were done under the temperature range varying from 25°C to 100°C. For each slot, the device performance was monitored and analyzed over the given temperature range. The results showed SiC DMOSFET had robust long-term reliability in high-power applications that are susceptible to pulse over currents, such as pulsed power modulators and hard-switched power electronics. The present devices on SiC boasts of higher breakdown, higher current handling capabilities, lower power dissipation and better thermal conductivity than those of silicon based power devices. Despite all the above superior properties, the device fabrication of Silicon Carbide MOS devices still has some distinct problems that hasn't been resolved completely yet, which limits the flexibility of SiC devices unlike Si based power devices. But with the rapidly maturing fabrication technologies this issue would be soon overcome.

# Chapter 3

## A COMPARATIVE STUDY ON UNIFORMLY DOPED AND LINEARLY GRADED PROFILE

### 3.1 Introduction

This chapter analyzes the two basic doping profiles for the drift region- uniformly doped profile and linearly graded doping profile respectively. The entire behavior of DIMOSFET (like other MOSFETs) revolve around the characteristics of p-n junction, which in this case is formed between the p-well and the  $n^-$  drift region. Due to the lower doping concentration at the drift region, the depletion region width extends entirely into the drift region and this charge free region sustains the reverse breakdown voltage for which the device is designed.

Here the analysis is made for a breakdown voltage as well as the power dissipation. A comparative discussion is done based on results present in [29]. As mentioned above, the base of device design is dependent on depletion width that extends into the lightly doped drift region which in turn is used to estimate the drift region thickness.

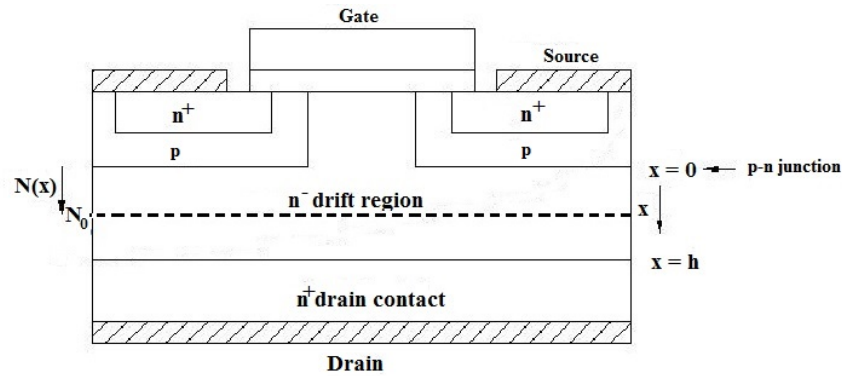
### 3.2 Vertical DIMOSFET with Uniformly Doped Profile

This type of profile has a uniformly doped concentration in the drift region (Region C) as shown in Fig.1.4. This is a special case of abrupt p-n junction when  $N_D \ll N_A$ , where  $N_A$  is the acceptor concentration of the p-body and  $N_D$  is the donor concentration of the drift region. The main objective to reduce the magnitude of  $R_{on-sp}$  is to reduce the power dissipation,  $P_D$  across

the vertical DIMOSFET, since[9]:

$$P_D \approx J_{on}^2 R_{on-sp} A \dots (3.1)$$

where  $J_{on}$  is the on-state current density and  $A$  is the cross-sectional area of the vertical device. Fig.3.1 illustrates the uniformly doped profile in the drift region of DIMOSFET. The peak doping concentration,  $N_o$  is constant and is same as  $N_{eff}$  for uniformly doped profile.



**Figure 3.1:** Uniformly Doped Profile in DIMOSFET

### 3.3 Vertical DIMOSFET with Linearly Graded Doping Profile

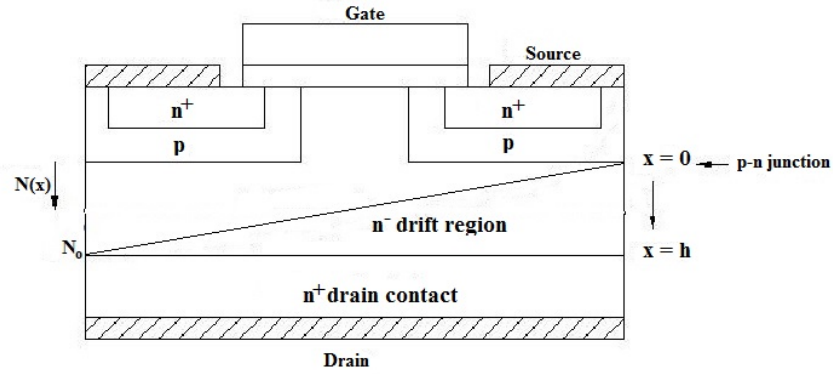
In practical devices, the doping profiles are not abrupt, especially near the metallurgical junction where the two types meet and compensate each other[35]. When depletion width terminate within the transition region, the doping profile can be approximated by a linear function. Unlike uniformly doped profile, this profile has a monotonically increasing profile along the three regions, namely region A, B and C(refer Fig.1.4). Such a profile has a lightly doped region at the top to a linearly increasing doping level(with  $N_o$  as the peak concentration value) towards the drain region at the bottom of the device as indicated in Fig.3.2.

The effective doping concentration( $N_{eff}$ ) is different, unlike in uniformly doped profile[29]. Here

$$N_{eff} = \frac{\alpha h}{\ln[1 + \frac{\alpha h}{N_o}]} \dots (3.2)$$

where  $\alpha$  is the concentration gradient (given by ratio of concentration difference to the height considered).

Using a linearly graded profile gives a wider depletion region due to low doping level near the



**Figure 3.2:** Linearly Graded Doping Profile in DIMOSFET

junction of p-body and  $n^-$  drift region and lower parasitic resistance due to the higher doping level towards the drain end, for the same peak concentration ( $N_o$ ).

### 3.4 Discussion of results

Based on the results stated in [29], a device height of  $73\mu\text{m}$  gave a  $V_{BAV}$  (Avalanche Breakdown Voltage) of  $6.4\text{kV}$  and  $V_{BPT}$  (Punch through Breakdown Voltage) of  $5\text{kV}$  in uniformly doped. But a linear doping profile gave  $V_{BAV}$  equal to  $6.74\text{kV}$  and  $V_{BPT}$  of  $6.75\text{kV}$  for  $\alpha$  value of  $1.58 \times 10^{18}\text{cm}^{-4}$  and device height of  $65\mu\text{m}$ . In the calculations in [29] the lower of the two voltages is considered as the breakdown voltage of the device. Thus, the former with uniformly doped profile had a breakdown voltage,  $V_{BPT}$ , of  $5\text{kV}$  while the latter with,  $V_{BAV}$ , equal to  $6.74\text{kV}$ . At  $h=82\mu\text{m}$ , using linear graded profile a maximum breakdown of  $8.75\text{kV}$  was achieved though designed for  $10\text{kV}$ . Compared to uniformly doping profile, linearly graded one has a wider depletion in region B (owing to low doping level) which supports higher breakdown voltages; moreover, the higher doping levels in region C (Fig.1.4) results in lower parasitic series resistance than the former. Lower parasitic resistance implies lower specific on-resistance ( $R_{on-sp}$ ) and consequently lower power dissipation,  $P_D$  as given by eq.3.1. This is observed in [29], as the maximum reduction in  $P_D$  that was achieved over the range of current densities ( $J$ ) considered ( $1-1000\text{ A/cm}^2$ ), was found to be  $61\%$  against uniformly doped drift region devices.

# Chapter 4

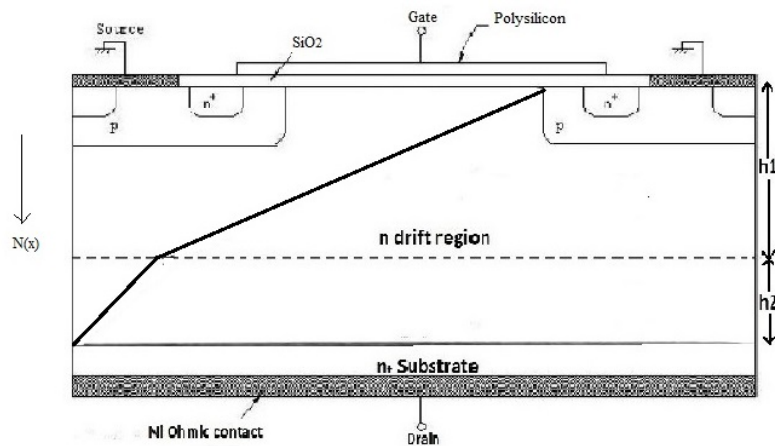
## ANALYSIS AND DESIGN OF 4H-DIMOSFET USING DISTORTED GAUSSIAN DOPING PROFILE

### 4.1 Introduction

This chapter analyzes structure of 4H-SiC DIMOSFET using distorted Gaussian doped profile in the drift region of the device. The device dimensions were determined by different parameters like the drift region doping level, blocking voltage, depletion region width, specific on-resistance and power dissipation. In the analysis carried out here, the device height was decided by iterations over the effective carrier concentration ( $N_{eff}$ ) and the distorted Gaussian profile was simplified to a linearly graded profile with two different concentration gradients ( $\alpha_1$  and  $\alpha_2$ ) along the entire device height. The slope of the profile (concentration gradient) was kept more steep along the lower end of the device to ensure lower values of parasitic resistances. Low doping ( $10^{12} - 10^{13}$  per cc) near the source end ensured higher breakdown for the device. The final effective carrier concentration ( $N_{eff}$ ) was obtained by iterations and the device height was calculated for each iteration based on  $N_{eff}$  value from the previous iteration. Separate calculations were done for effective doping concentration ( $N_{eff}$ ) and depletion region width. This was followed by the evaluation of the drain as well as channel voltages, specific on-resistance, power dissipation and breakdown voltages for current densities varying from 1-1000 A/cm<sup>2</sup>. The results so obtained have been analyzed by considering variations of current density versus power dissipation for different values of peak carrier concentration ( $N_o$ ), current density versus power dissipation for different  $N_o$ , punch through breakdown voltage versus avalanche breakdown voltage, avalanche breakdown voltage versus critical electric field and  $N_o$  versus concentration gradient. As compared to uniformly doped drift region the linearly graded profile shows a 61% reduction in power dissipation with the latter being thinner device [28-30].

## 4.2 Vertical DIMOSFET with Distorted Gaussian Profile in the Drift Region

It has already been mentioned earlier that there exists a trade-off between breakdown voltage and the specific on-resistance of drift region. Ideally, for any high power application we would want a lower value of parasitic resistance and at the same time a high reverse breakdown voltage. The trade-off exists because a lower resistance value requires higher doping levels and a smaller device height while a higher breakdown voltage requires low doping with a larger height of device. The main goal in reducing the  $R_{on-sp}$  of DIMOSFET is to reduce its power dissipation ( $P_D$ ) which is given by eq.3.1. The objective suggested in this chapter has been achieved by using a dual slope linearly graded profile in the epitaxial layer of the device with a lightly doped region near the p-well-n<sup>-</sup> junction to a linearly increasing doping level near the lower end of device. Such a profile is shown in Fig.4.1 below. Referring to the earlier discussion on Fig.1.4 the profile has a linearly graded profile in region A (accumulation-region), region B (JFET-region) with a steeper slope in region C (drift-region) of the DIMOSFET. Such a device structure should give a wide depletion region width due to low doping level in JFET region and low parasitic resistance in region C owing to higher level of dopant concentration here.



**Figure 4.1:** Distorted Gaussian Doping Profile in DIMOSFET

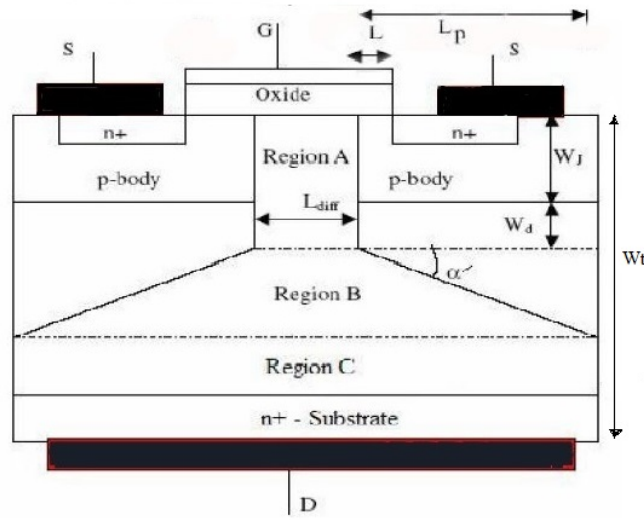
## 4.3 Theoretical Analysis

The basic structure of the DIMOSFET for the analysis with the device dimensions and the respective symbols is shown in Fig.4.2. This section discusses the device with a distorted Gaussian profile in the drift region which is simplified by a dual slope linearly graded profile. The formula for power dissipation,  $P_D$  as given by eq.3.1 was just an approximate value in which the effects of leakage current density ( $J_L$ ) was ignored. For a 50% duty cycle for various current

levels, the basic equation is given by[30]:

$$P_D = \frac{1}{2}(J_{on}^2 AR_{on-sp} + J_L V_B A) \dots (4.1)$$

where  $J_{on}$  is the on-state current density,  $J_L$  is the reverse leakage current density of junction between the p-well/n<sup>-</sup> drift region,  $V_B$  is the breakdown voltage of the device and  $A$  is the area of the device. Equation 4.1 is simplified to eq.3.1 since the value of  $J_L$  is negligible compared to the on-state current density(typically in  $\mu A/cm^2$ ).



**Figure 4.2:** Basic Device Structure of DIMOSFET[30]

The drain current( $I_{DS}$ ) which equals the channel current( $I_{CH}$ ) is given by[24]:

$$I_{DS} = \frac{W\mu_n}{2L[1 + (\mu_n/2v_{sat}L)V_{ch}]} V_{ch}[2C_{ox}(V_{GS} - V_T) - (C_{ox} + C_{do})V_{ch}] \dots (4.2)$$

where  $W$  is the width of the device, channel length is given by  $L$ ,  $\mu_n$  is the effective zero field doping dependent carrier mobility related to the doping level of the drift region,  $V_{ch}$  is the channel voltage,  $v_{sat}$  is the saturated carrier drift velocity with a value of  $2 \times 10^7$  cm per sec,  $C_{ox}$  is the oxide capacitance per unit area,  $V_{GS}$  is the gate-source voltage and  $C_{do}$  is the depletion capacitance of the body.  $C_{do}$  is neglected as it's value is negligible compared to  $C_{ox}$ . In the equation above, the effect of field on mobility is considered while formulating the value of drain current. The voltage drops across the three regions, viz. region A,B and C have been found as [24]:

$$V_A = \frac{I_{DS}(W_j + W_d)}{W(L_{diff}qN_{eff})\mu_n - I_{DS}/E_c} \dots (4.3)$$

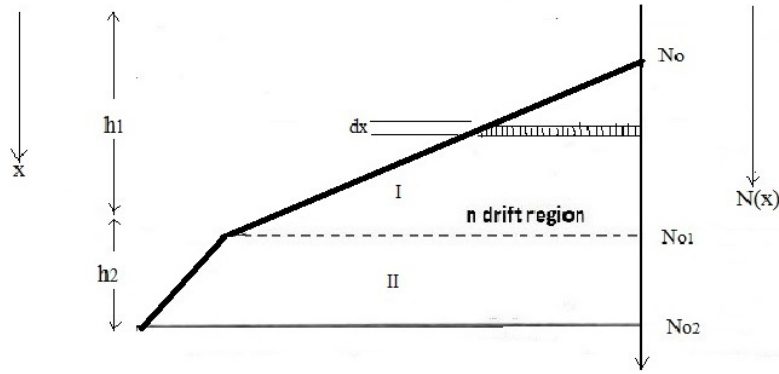
$$V_B = \frac{I_{DS}}{WqN_{eff}\mu_n \cot\alpha} \log\left[\frac{WqN_{eff}\mu_n(L_{diff} + L_p) - I_{DS}/E_c}{WqN_{eff}\mu_n L_{diff} - I_{DS}/E_c}\right] \dots (4.4)$$

$$V_C = \frac{I_{DS}(W_t - W_j - W_d - L_p \tan\alpha)}{WqN_{eff}\mu_n(L_{diff} + 2L_p) - I_{DS}/E_c} \dots (4.5)$$

where all the parameters are as shown in Fig.4.2.  $L_{diff}$  is the p-body separation,  $W_t$  is the device height which is set by using the depletion width at punch through breakdown,  $W_j$  is the p-well thickness and  $W_d$  is the on-state depletion region width.  $V_A$ ,  $V_B$  and  $V_C$  are the voltage drops across regions A, B and C respectively. The drain to source voltage ( $V_{DS}$ ) is obtained by adding  $V_A$ ,  $V_B$  and  $V_C$  and  $V_{ch}$ , i.e,  $V_{DS} = V_A + V_B + V_C + V_{ch}$ .

### 4.3.1 Calculation of effective carrier concentration( $N_{eff}$ )

To calculate the  $N_{eff}$  we use the doping profile graph as shown in Fig.4.3. The graph shows the variation of doping concentration along the device height. The total effective concentration ( $N_{eff}$ ) can be considered as the summation of effective concentrations over regions I and II.



**Figure 4.3:** Distorted Gaussian Profile

For  $0 \geq x \leq h_1$  :

$$N(x) = N_o + \alpha_1 x \dots (4.6)$$

for  $h_1 \geq x \leq h_2$

$$N(x) = N_{o1} + \alpha_2 x \dots (4.7)$$

where  $\alpha_1$  and  $\alpha_2$  are the concentration gradients of regions I and II respectively. The resistance of region I ( $R_I$ ) is calculated by considering an elemental horizontal strip of width  $dx$  and integrating over height  $h_1$ .

$$R_I = \int_0^{h_1} \frac{dx}{A\mu_n q(N_o + \alpha_1 x)} \dots (4.8)$$

Integrating the above equation gives the value of resistance as (refer Appendix for derivation)

$$R_I = \frac{1}{\mu_n q A \alpha_1} \ln \left[ 1 + \frac{\alpha_1 h_1}{N_o} \right] \dots (4.9)$$

Equating the above result with the general formula of resistance,  $R = h_1 / N_{eff} q \mu_n A$  gives the value of effective concentration over region I as:

$$N_{eff1} = \frac{\alpha_1 h_1}{\ln \left[ 1 + \frac{\alpha_1 h_1}{N_o} \right]} \dots (4.10)$$

Similarly, for region II the resistance value,  $R_{II}$  is calculated by integrating as below:

$$R_{II} = \int_{h_1}^{h_1+h_2} \frac{dx}{A\mu_n q(N_{o1} + \alpha_2 x)} \dots (4.11)$$

As above equating the obtained value of resistance with the general formula of resistance gives the effective doping concentration of region II ( $N_{eff2}$ ) as:

$$N_{eff2} = \frac{\alpha_2 h_2}{\ln \left[ 1 + \frac{\alpha_2 h_2}{N_{o1} + \alpha_2 h_1} \right]} \dots (4.12)$$

The total effective carrier concentration ( $N_{eff}$ ) in the drift region is obtained by summing eq.4.11 and 4.12

$$N_{eff} = \frac{\alpha_1 h_1}{\ln \left[ 1 + \frac{\alpha_1 h_1}{N_o} \right]} + \frac{\alpha_2 h_2}{\ln \left[ 1 + \frac{\alpha_2 h_2}{N_{o1} + \alpha_2 h_1} \right]} \dots (4.13)$$

For 4H and 6H SiC, the carrier mobility's dependence on doping concentration is given by[4]:

$$\mu_n = 1.93 \times 10^8 N_{eff}^{-0.34} \text{ cm}^2/\text{Vs} \dots (4.14)$$

However, these experimental datas are valid only for  $N_{eff}$  greater than  $3 \times 10^{16}$ , for values less than  $3 \times 10^{16}$ , the peak value of  $530 \text{ cm}^2/\text{Vs}$  is considered.

### 4.3.2 Calculation of Channel Voltage ( $V_{ch}$ )

$V_{ch}$  is evaluated by using eq.4.2. In this equation we had discussed that the value of  $C_{do}$  is much less compared to  $C_{ox}$  and hence can be neglected. Keeping  $V_{GS} = 40\text{V}$  and  $V_T = 1\text{V}$ , eq.4.2 becomes[30]

$$I_{ch} = \frac{W\mu_n C_{ox} V_{ch} [78 - V_{ch}]}{2L[1 + (\mu_n/2v_{sat}L)V_{ch}]} \dots (4.14)$$

On simplifying the above equation we get

$$I_{ch} = \frac{W\mu_n C_{ox} V_{ch} [78 - V_{ch}]}{2v_{sat}L + \mu_n V_{ch}} \dots (4.15)$$

eq.4.15 can again be written as

$$W\mu_n v_{sat} C_{ox} V_{ch}^2 + (\mu_n I_{ch} - 78W\mu_n v_{sat} C_{ox}) V_{ch} + 2v_{sat} L I_{ch} = 0 \dots (4.16)$$

eq.(4.16) being a quadratic equation, the value of  $V_{ch}$  can be evaluated as:

$$V_{ch} = \frac{(\mu_n I_{ch} - 78W\mu_n v_{sat} C_{ox}) \pm \sqrt{(\mu_n I_{ch} - 78W\mu_n v_{sat} C_{ox})^2 - 8LW I_{ch} \mu_n v_{sat}^2 C_{ox}}}{2W\mu_n v_{sat} C_{ox}} \dots (4.17)$$

The last parameter to be evaluated is the specific on-resistance ( $R_{on-sp}$ ) of the DIMOSFET, which is defined as the product of device area (A) and the resistance of drift region (R).  $R_{on-sp}$  is expressed as[30]:

$$R_{on-sp} = RA = \rho.l = \frac{(W_t - W_j - W_d - L_p \tan \alpha')}{\mu_{eff} q N_{eff}} \dots (4.18)$$

where  $\alpha'$  is the angle of the slope of the drift region narrowing and the value of  $\mu_{eff}$  is obtained from eq.(4.14) corresponding to the effective concentration ( $N_{eff}$ ) of the linearly graded drift region.

### 4.3.3 Critical Field, $E_c$ and Breakdown Voltage Calculations

The two device breakdown mechanisms considered here is - Punch through and Avalanche Breakdown. The lower of these two breakdown voltages is considered as the breakdown voltage of the DIMOSFET. The depletion width at breakdown,  $W$  is first estimated by estimating the width at punch through voltage,  $V_{BPT}$  obtained by[35]:

$$W = \left[ \frac{12\epsilon_s(V_{BPT} + V_g)}{qN_{eff}} \right] \dots (4.19)$$

where  $\epsilon_s$  is the permittivity value for 4H-Silicon Carbide equal to  $9.66\epsilon_o$ . Now as  $V_g \ll V_{BPT}$ , where  $V_g$  is the gradient voltage for linearly graded profile in the drift region. This assumption reduces eq.(4.19) to:

$$W = \left[ \frac{12\epsilon_s V_{BPT}}{qN_{eff}} \right] \dots (4.20)$$

The critical field,  $E_c$  corresponding to the depletion region width( $W$ ) obtained as above could be obtained by using the equation[35]:

$$E_c = \left( \frac{q\alpha W^2}{8\epsilon_s} \right) \dots (4.21)$$

The avalanche breakdown voltage ( $V_{BAV}$ ) can then be calculated using the equation[35]:

$$V_{BAV} = \frac{2}{3} E_c W \dots (4.22)$$

where the depletion width at the two breakdown voltages is set equal to each other. To begin with we set the  $V_{BPT}$ , calculate the depletion width  $W$  corresponding to the punch through breakdown voltage. The value of  $W$  is then used to calculate the critical field  $E_c$ ,  $V_{BAV}$  is finally calculated by using eq.(4.21). The minimum of the two breakdown voltages' values is considered as the breakdown voltage of the device.

## 4.4 Calculations and Related Graphs

The dimensions of 4H DIMOSFET has been set so that height of region I,  $h_1$  equals the depletion region width,  $W$  under a reverse bias voltage of 10kV applied on the p-n junction formed between the p-body/ $n^-$ -epitaxial layer. The DIMOSFET was thus, designed for a maximum blocking voltage of 10kV. The dimensions of other variables labeled in Fig.4.2 have been taken as :  $W_t=1\mu\text{m}$ ,  $L_p=25\mu\text{m}$ . The value of  $W_t$  is taken  $1\mu\text{m}$  as the implant depth in 4H and 6H-SiC is of this order[29]. The cross-sectional area of the device is  $24000 \times 10^{-8} \text{ cm}^2$  ( $300\mu\text{m} \times 80\mu\text{m}$ ) but the analysis is carried out for only one half of the device due to symmetrical device structure, i.e,  $A=12000 \times 10^{-8} \text{ cm}^2$ .

Calculations of various parameters of the 4H-SiC DIMOSFET with a distorted Gaussian drift region is made by using a doping level of  $10^{12}$  per cc near the source end to  $10^{15}$  per cc over height  $h_1$  and  $10^{15}$  per cc to  $10^{16}$  per cc near the drain end in region II. The device heights( $h_1$  and  $h_2$ ) has been set by iterations over the effective carrier concentration. Initially  $N_{eff1}$  (and  $N_{eff2}$ ) is chosen as the mid value and the corresponding value of  $h_1$ (and  $h_2$ ) is calculated. This value of  $h_1$  is then used to calculate  $N_{eff1}$ (and  $N_{eff2}$ ) iteratively by using eq.(4.10) and eq.(4.12). The value of slope gradients ( $\alpha_1$  and  $\alpha_2$ ) is obtained by taking difference in carrier concentration of the two extreme values in region I and II respectively. The above procedure is repeated for values  $N_o$  equal to  $2 \times 10^{12}$ ,  $2.5 \times 10^{12}$  and  $10^{13}$  per cc. The specific on-resistance ( $R_{on-sp}$ ) is then obtained from the calculated value of  $N_{eff}$  as per eq.(4.18). For region I,  $R_{on-sp1}$  is calculated from  $N_{eff1}$  and  $R_{on-sp2}$  for region II is calculated from  $N_{eff2}$  by using the same equation eq.(4.18). Moreover, the value of  $\mu_{eff}$  to be used in eq.(4.18) is obtained from eq.(4.14).

The values of power dissipation,  $P_D$  are then calculated for different current densities,  $J_f$  (same as  $J_{on}$ ) ranging from 1-1000 A/cm<sup>2</sup>. This is then repeated for different concentration gradients ( $\alpha_1$  and  $\alpha_2$ ). The results of the calculations are shown in Tables 4.1 to 4.4. This is followed by plots of current density versus power dissipation for different values of  $N_o$ , current density versus on-state depletion width( $W_d$ ) for different values of  $N_o$ . The variation of concentration gradient( $\alpha_1$ ) for various values of  $N_o$  is shown in Fig.(4.7). The variations of avalanche breakdown voltage( $V_{BAV}$ ) versus critical field ( $E_c$ ) and punch through breakdown voltage ( $V_{BPT}$ ) and avalanche breakdown voltage for various values of  $N_o$  are shown in Fig.(4.8) and Fig.(4.10) respectively. Fig.(4.11) illustrates the variation of on-state depletion width to specific on-resistance, the four successive nodal points on each plot corresponds to the values of current density varying from 1 A/cm<sup>2</sup> to 1000 A/cm<sup>2</sup>.

**Table 4.1: Results of currents, voltages,  $R_{on-sp}$  and  $P_D$  for doping  $10^{12}$ - $10^{15}$  in region I and  $10^{15}$ - $10^{16}$  per cc in region II,  $h_1=0.0272\text{cm}$ ,  $N_{eff}=2.85 \times 10^{17}/\text{cc}$ ,  $\alpha_1=3.674 \times 10^{16} \text{ cm}^{-4}$**

J(A/cm <sup>2</sup> )	$W_d(\text{cm})$	$E_c(\text{V/cm})$	$R_{on-sp1}(\Omega - \text{cm}^2)$	$P_D(\text{W})$
1	5.5e-7	6.35e5	0.0038	2.25e-7
10	1.74e-6	6.35e5	0.0035	2.12e-5
100	5.54e-6	6.35e5	0.0028	0.0017
1000	1.88e-5	6.35e5	0.00022	0.0136

**Table 4.2: Results of currents, voltages,  $R_{on-sp}$  and  $P_D$  for doping  $2 \times 10^{12}$ - $10^{15}/\text{cc}$  in region I and  $10^{15}$ - $10^{16}/\text{cc}$  in region II,  $h_1=0.0258\text{cm}$ ,  $N_{eff}=2.65 \times 10^{17}/\text{cc}$ ,  $\alpha_1=3.87 \times 10^{16} \text{ cm}^{-4}$**

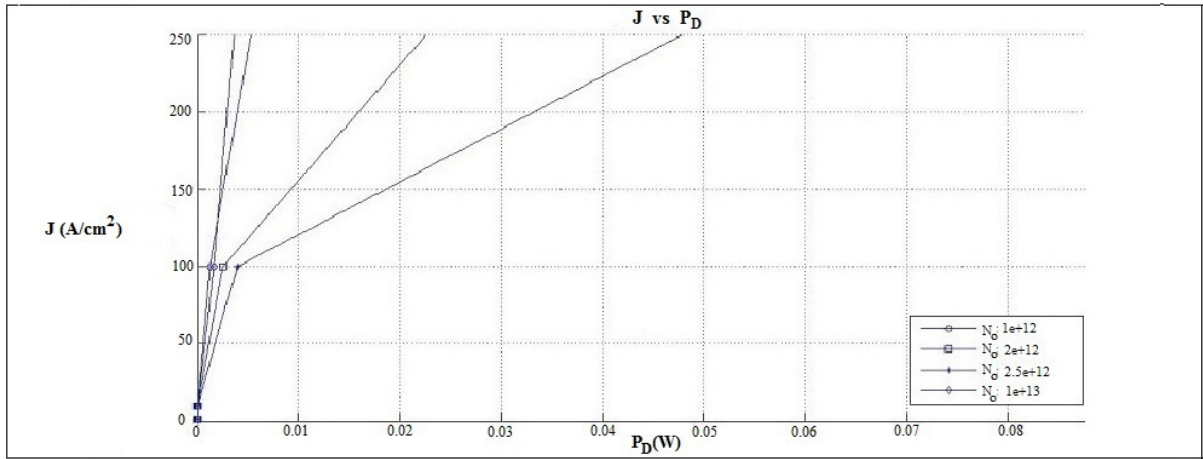
J(A/cm <sup>2</sup> )	$W_d(\text{cm})$	$E_c(\text{V/cm})$	$R_{on-sp1}(\Omega - \text{cm}^2)$	$P_D(\text{W})$
1	5.65e-7	6.021e5	0.0052	3.13e-7
10	1.79e-6	6.021e5	0.0050	3.001e-5
100	5.69e-6	6.021e5	0.0044	0.0026
1000	1.93e-5	6.021e5	0.00020	0.1226

**Table 4.3: Results of currents, voltages,  $R_{on-sp}$  and  $P_D$  for doping  $2.5 \times 10^{12}$ - $10^{15}/\text{cc}$  in region I and  $10^{15}$ - $10^{16}/\text{cc}$  in region II,  $h_1=0.0253\text{cm}$ ,  $N_{eff}=2.58 \times 10^{17}/\text{cc}$ ,  $\alpha_1=3.94 \times 10^{16} \text{ cm}^{-4}$**

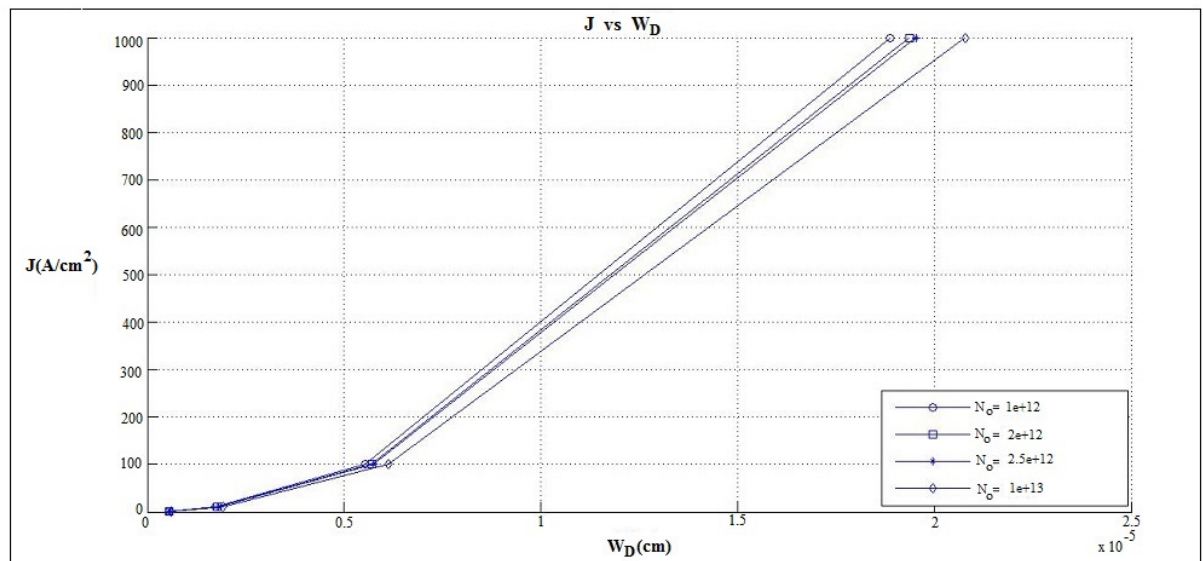
J(A/cm <sup>2</sup> )	$W_d(\text{cm})$	$E_c(\text{V/cm})$	$R_{on-sp1}(\Omega - \text{cm}^2)$	$P_D(\text{W})$
1	5.7e-7	5.91e5	0.0075	4.51e-7
10	1.8e-6	5.91e5	0.0073	4.39e-5
100	5.74e-6	5.91e5	0.0067	0.0040
1000	1.95e-5	5.91e5	0.00045	0.2671

**Table 4.4: Results of currents, voltages,  $R_{on-sp}$  and  $P_D$  for doping  $10^{13}$ - $10^{15}/cc$  in region I and  $10^{15}$ - $10^{16}/cc$  in region II,  $h_1=0.0223cm$ ,  $N_{eff}=2.15 \times 10^{17}/cc$ ,  $\alpha_1=4.44 \times 10^{16} cm^{-4}$**

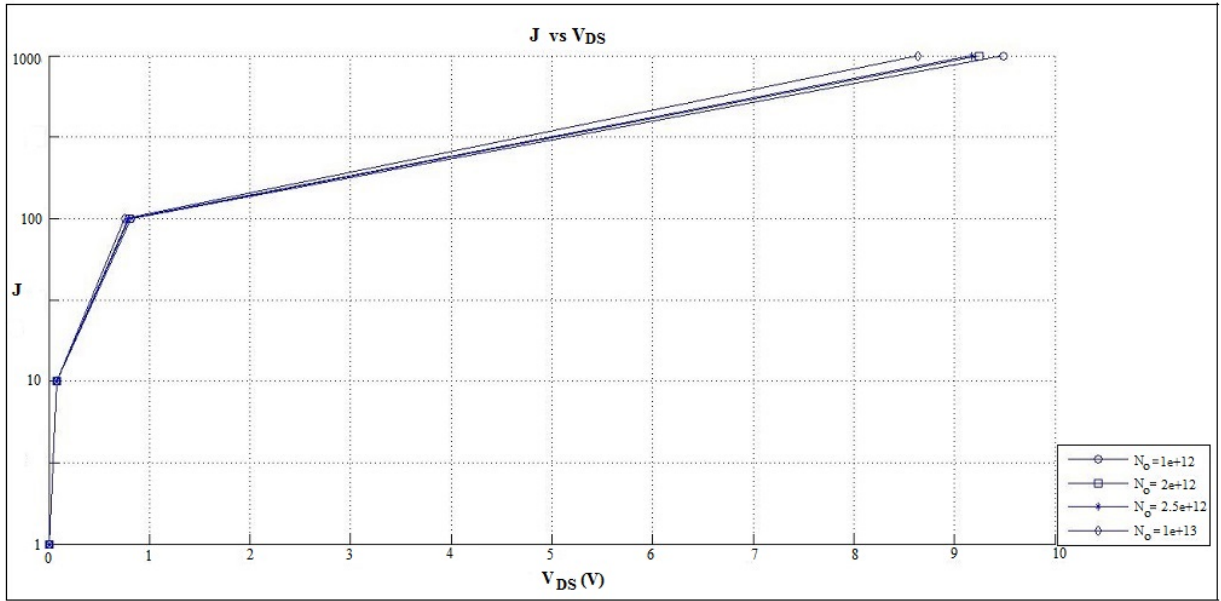
$J(A/cm^2)$	$W_d(cm)$	$E_c(V/cm)$	$R_{on-sp1}(\Omega - cm^2)$	$P_D(W)$
1	6.09e-7	5.16e5	0.0028	1.67e-7
10	1.93e-6	5.16e5	0.0026	1.68e-5
100	6.13e-6	5.16e5	0.0021	0.0013
1000	2.08e-5	5.16e5	0.00042	0.0255



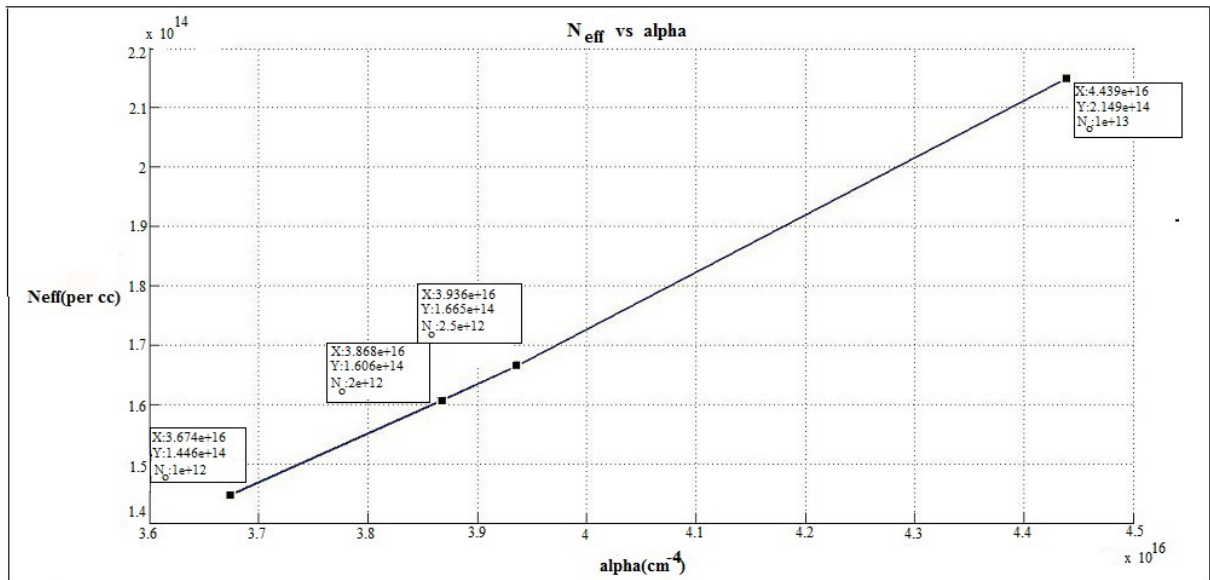
**Figure 4.4:** Plot of current density vs power dissipation for 4H-SiC DIMOSFET for different values of  $N_o$



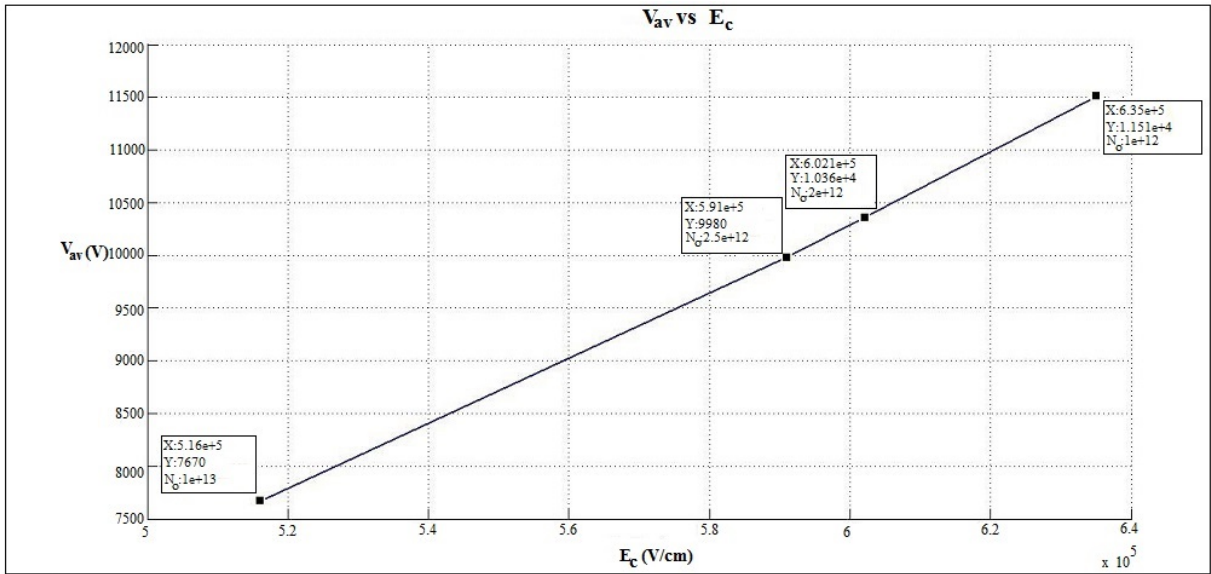
**Figure 4.5:** Plot of current density vs on-state depletion width for 4H-SiC DIMOSFET for different values of  $N_o$



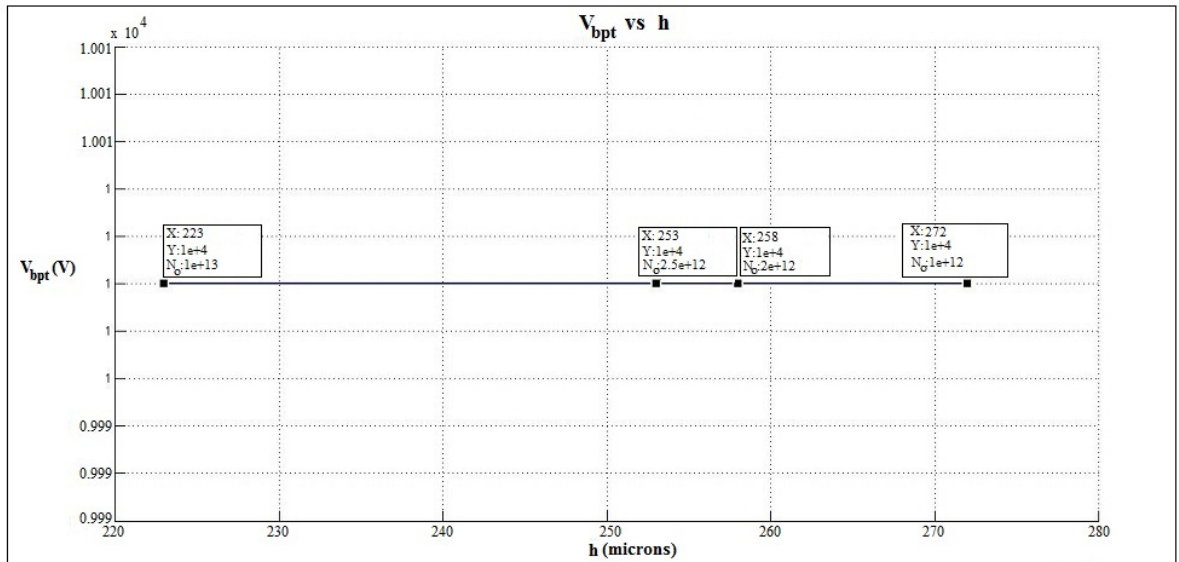
**Figure 4.6:** Plot of current density vs drain-to-source voltage for 4H-SiC DIMOSFET for different values of  $N_0$



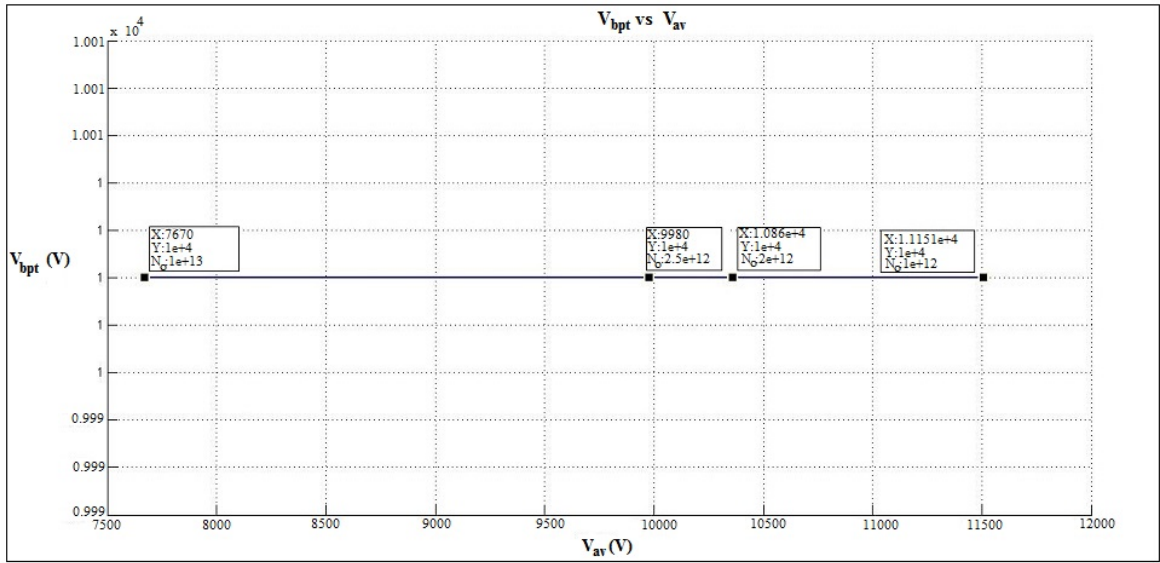
**Figure 4.7:** Plot of effective doping concentration vs concentration gradient(region I) for 4H-SiC DIMOSFET



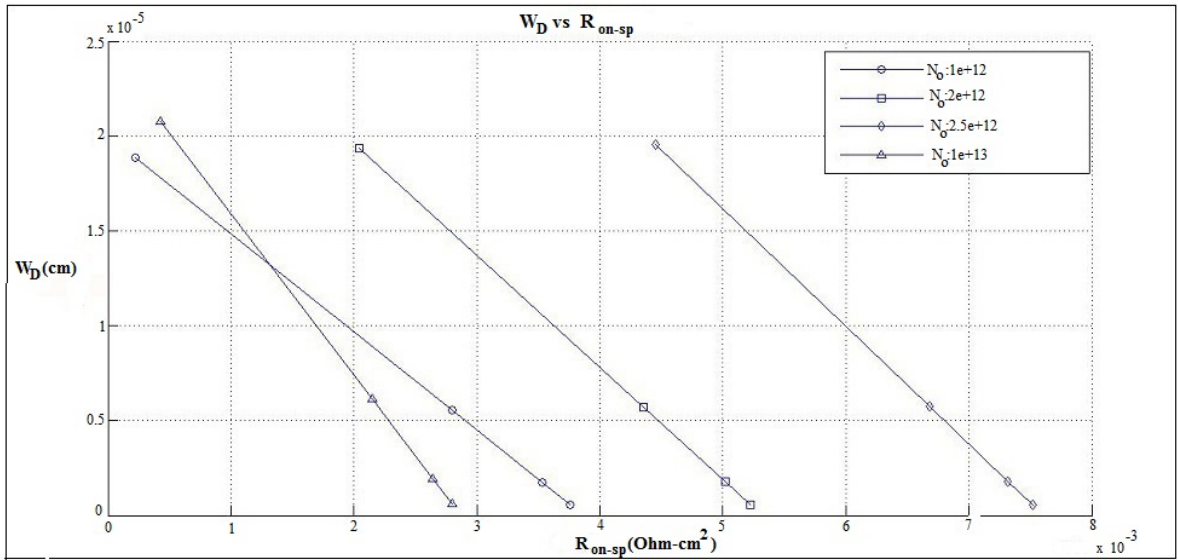
**Figure 4.8:** Plot of avalanche breakdown voltage vs critical field for 4H-SiC DIMOSFET for different values of  $N_o$



**Figure 4.9:** Plot of punch through breakdown voltage vs device height(region I) for different values of  $N_o$



**Figure 4.10:** Plot of punch through breakdown voltage vs avalanche breakdown voltage for different values of  $N_o$



**Figure 4.11:** Plot of on-state depletion width vs specific-on resistance voltage for different values of  $N_o$

#### 4.4.1 Calculation of Breakdown Voltages( $V_{BAV}$ and $V_{BPT}$ )

As indicated in sub section(4.3.3), two breakdown voltages are considered here the avalanche breakdown voltage( $V_{BAV}$ ) and punch through voltage( $V_{BPT}$ ). The device design requires the two voltages to be made equal to each other(or almost equal) for a given voltage. Here, the lower of the two is considered as the breakdown voltage of the DIMOSFET.

For the doping profile considered, the punch through breakdown voltage is set to 10kV and the device height  $h_1$ , (since  $h_2 \ll h_1$ ) is set equal to the depletion region width. The value of critical field,  $E_c$  is calculated using eq.(4.21), the value of  $E_c$  so obtained is then used to calculate  $V_{BAV}$  from equation eq.(4.22). The results obtained is shown in table(4.5) below:

**Table 4.5: Results of breakdown voltages,  $V_{BAV}$  for a punch through breakdown,  $V_{BPT}$  set to 10kV**

$\alpha_1(\text{cm}^{-4})$	$E_c(\text{V/cm})$	$V_{BAV}(\text{V})$
$3.67 \times 10^{16}$	$6.35 \times 10^5$	11.51
$3.87 \times 10^{16}$	$6.02 \times 10^5$	10.36
$3.94 \times 10^{16}$	$5.91 \times 10^5$	9.98
$4.44 \times 10^{16}$	$5.16 \times 10^5$	7.67

#### 4.5 Comparison between uniformly doping profile and distorted Gaussian profile

As discussed in Section(3.4), the maximum breakdown voltage for which uniformly doping can be used is 5kV. Here we use the same device dimensions and doping concentrations (peak doping concentration for distorted profile) to compare the results of the two above mentioned doping profiles. Table(4.7) lists out the different parameters for uniformly doping profile and Table(4.6) lists the same set of parameters for distorted profile. For the comparison of breakdown voltages, a punch through breakdown voltage,  $V_{BPT}$  of 5kV is set for both the profiles. Table(4.8) compares the avalanche breakdown voltage( $V_{BAV}$ ), critical electric field( $E_c$ ), device height and carrier concentrations for both profiles.

**Table 4.6: Results of currents,voltages,  $R_{on-sp}$  and  $P_D$  for doping  $10^{15}/cc$ ,  $h=0.0073cm$  for uniformly doped drift region**

J(A/cm <sup>2</sup> )	$W_d(cm)$	$E_c(V/cm)$	$R_{on-sp}(\Omega - cm^2)$	$P_D(W)$
1	2.61e-5	1.78e6	0.0359	2.16e-6
10	8.3e-5	1.78e6	0.0331	2.11e-4
100	2.67e-4	1.78e6	0.0292	0.0199
1000	8.82e-4	1.78e6	0.024	1.552

**Table 4.7: Results of currents,voltages,  $R_{on-sp1}$  and  $P_D$  for doping  $10^{15}$ - $10^{19}/cc$  in region I and  $10^{19}$ - $10^{20}/cc$  in region II,  $h_1=0.0158cm$ , $N_{eff}=1.09 \times 10^{18}/cc$ ,  $\alpha_1=1.54 \times 10^{21} cm^{-4}$  for distorted profile**

J(A/cm <sup>2</sup> )	$W_d(cm)$	$E_c(V/cm)$	$R_{on-sp1}(\Omega - cm^2)$	$P_D(W)$
1	7.77e-6	4.20e6	9.32e-5	5.59e-7
10	1.67e-5	4.20e6	8.68e-5	5.55e-5
100	3.58e-5	4.20e6	7.65e-5	5.52e-5
1000	8.46e-5	4.20e6	6.79e-5	0.0055

**Table 4.8: Results of calculation of breakdown voltages, $V_{BAV}$  for a set punch through voltage,  $V_{BPT} =5kV$  for uniformly doped profile and distorted Gaussian profile**

Parameters	Uniformly-doped	Distorted Gaussian
$V_{BAV}$	6.4 kV	5.2 kV
$E_c(V/cm)$	$1.78 \times 10^6$	$1.86 \times 10^6$
Device height, $h(\mu m)$	73	42
$\alpha(cm^{-4})$	–	$4.5 \times 10^{20}$
$N_{eff}/(cc)$	$10^{15}$	$5.5 \times 10^{15}$

## 4.6 Discussion of Results

Tables 4.1 - 4.4 shows the variations of on-state depletion width  $W_d$ , critical electric field  $E_c$ , specific on-resistance  $R_{on-sp}$  and power dissipation  $P_D$  for peak impurity concentration  $N_o$  values  $10^{12}$ ,  $2 \times 10^{12}$ ,  $2.5 \times 10^{12}$  and  $10^{13}$  per cc over the range of current densities considered i.e, 1-1000 A/cm<sup>2</sup>. These four tables also indicate that the value of  $N_{eff}$  increases in an almost linear fashion with  $\alpha 1$ , as is evident in eq.4.10. The variations of concentration gradient( $\alpha 1$ ) over region I, critical electric field ( $E_c$ ) and avalanche breakdown voltage ( $V_{BAV}$ ) with peak impurity concentration ( $N_o$ ) is shown in table 4.5. Since critical field  $E_c$  decreases with lower avalanche breakdown voltage, the values of increasing  $N_o$  reduces the value of  $\alpha 1$  which also reduces the value of  $E_c$ .  $V_{BAV}$  is seen varying linearly with changing  $E_c$  which is as per eq.4.22. Table 4.8 compares the different parameters of uniformly doped and distorted Gaussian doping profiles for a set punch through breakdown voltage of 5kV. The device height for latter profile is almost half of that of former for the same breakdown voltage. The avalanche breakdown voltage then calculated for uniformly doped profile is 6.4kV for device height of  $73\mu m$  and 5.2kV for distorted Gaussian profile with a device height of  $42\mu m$ . Similar to the discussion in section 3.4, for a device design both the above breakdown voltages( $V_{BPT}$  and  $V_{BAV}$ ) are needed to be made equal to each other and as in calculations in [29], the lower of the two voltages is considered as the breakdown voltage of the device.

## Chapter 5

# CONCLUSION AND FUTURE SCOPE OF WORK

The results of the calculations based on the proposed model using two step linearly graded profile, which is an approximation of a Gaussian profile in the drift region of a 4H- DIMOSFET have been presented in Chapter 4 of this work and have been quoted graphically through figures Fig.4.4 to Fig.4.11. The main variables that have been used are namely, the current density  $J$ , the depletion region width  $W_d$  in the drift region, the drain-source voltage  $V_{DS}$ , the effective concentration of impurities  $N_{eff}$  in the drift region, the critical electric field  $E_c$ , the avalanche breakdown voltage  $V_{BAV}$ , the punch through breakdown voltage  $V_{BAV}$ , the power dissipation  $P_D$ , the device height  $h$  and also the specific on-resistance  $R_{on-sp}$  of the device, which equals to a first approximation to the  $R_{on-sp}$  of the drift region for high voltage devices.

The variations of current density versus power dissipation of the device treating the peak concentration  $N_o$  as a parameter have been shown in Fig.4.4. It has been observed that the  $P_D$  rises steeply from low values of current density to a value of  $J$  of about  $1000 \text{ A/cm}^2$ . The values of  $P_D$  continue to rise steeply with  $J$  for  $N_o$  equal to  $1 \times 10^{13}$  and  $1 \times 10^{12}$  per cc. There is however, a crossover at very low values of  $P_D$  between the two graphs close to the origin . For the other two graphs, namely for  $N_o$  equal to  $2 \times 10^{12}$  and  $2.5 \times 10^{12}$  per cc, the  $P_D$  versus  $J$  graph has an inflexion point much less than  $0.01 \text{ W}$ . The former concentration rises more steeply than the latter with increasing values of  $J$ .

The current density  $J$  versus depletion region width  $W_d$  shows a gradual rise in  $W_d$  as the value of  $N_o$  rises from  $N_o = 1 \times 10^{12}$  to  $2 \times 10^{12}$ . The similar profiles are observed for higher values of  $W_d$  versus  $J$  for  $N_o = 2.5 \times 10^{12}$  and  $1 \times 10^{13}$  per cc as shown in Fig.4.5. Here again there

lies an inflexion point at values of  $J = 100 \text{ A/cm}^2$  with  $W_d$  slightly greater than 0.05 microns. The plot of current density  $J$  versus  $V_{DS}$  in Fig.4.6 shows similar nature of variation for values of  $N_o$  ranging from  $1 \times 10^{12}$  to  $1 \times 10^{13}$  per cc. This graph also shows an inflexion point close to the drain to source voltage  $V_{DS}$  less than 1V.

The graph of  $N_{eff}$  versus concentration gradient  $\alpha$  has a linear plot with the lowest value at  $N_o = 1 \times 10^{12}/\text{cc}$  to  $N_o = 1 \times 10^{13}/\text{cc}$ . No such inflexion point has been observed in this graph as has been quoted earlier. This is quite evident from equation eq.4.10 wherein the value of  $N_{eff}$  increases in an almost linear fashion with  $\alpha$  as the denominator. In that equation containing  $\alpha$  doesn't significantly affect the value of  $N_{eff}$ . The variations of critical field  $E_c$  with the avalanche breakdown voltage has been plotted in Fig.4.8 with  $N_o$  as a parameter. Since the critical field  $E_c$  decreases with lower avalanche breakdown voltage, the values of increasing  $N_o$  reduces the value of  $\alpha$  which also reduces the value of  $E_c$ .

The device height  $h$  was pre-calculated on the basis of the depletion region width obtained by the normal device equation. Hence, an increase in the value of  $N_o$  lowers the value of  $\alpha$  and hence the punch through breakdown voltage giving a lower value of depletion region width which had been set equal to the device height  $h$ . This shows the nature of variation of punch through voltage with  $h$  as shown in Fig.4.9. Along with this argument it can be said that a 10kV DIMOSFET using the Gaussian profile approximated to a linearly graded profile will have an increasing value of  $h$  with decreasing value of  $N_o$ . The variation in values of punch through breakdown voltage  $V_{BPT}$  with the avalanche breakdown voltage  $V_{BAV}$  is shown in Fig.4.10. It can be seen that the two voltages are equal at the pre-designed breakdown voltage of 10kV for which  $N_o = 2.5 \times 10^{12}/\text{cc}$ . This is the ideal device using the specified Gaussian profile in the drift region of the 4H-SiC DIMOSFET, as the best device design is the one in which these two breakdown voltages equal each other. For other points on this graph, the design can't be approved as the breakdown voltage having a lower value will come into force before the other breakdown voltage can be reached.

The variation of  $W_d$  with the specific-on resistance  $R_{on-sp}$  for various values of  $N_o$  are shown in Fig.4.11. It can be seen that the magnitude of  $R_{on-sp}$  decreases as the value of  $W_d$  increases. This can be explained with reference to eq.4.18 wherein the depletion region width has a negative slope with respect to  $R_{on-sp}$ . The value of  $N_{eff}$  depended upon  $N_o$  has been used in that equation and it can be seen as the device height  $h (\approx W_t)$  the numerator in eq.4.18 decreases as  $W_d$  increases. The nature of variation shown in Fig.4.4 of  $J$  versus  $P_D$  wherein the cross over

has been observed for values of  $N_o=10^{12}$  and  $10^{13}$  per cc the graph of Fig.4.11 shows a similar crossover for these values of  $N_o$ .

In order to conclude, it can be said that the best device using this type of profile is the one which can be obtained from Fig.4.10 with  $N_o= 2.5 \times 10^{12}$  per cc,  $\alpha= 3.939 \times 10^{16} \text{cm}^{-4}$  (refer Fig.4.7) having  $V_{BPT} = V_{BAV} = 10\text{kV}$ . The crossover points shown in Fig.4.4 and Fig.4.11 couldn't be explained on the basis of current analysis. This can be undertaken by other researchers for future work on these devices having different polytypes of SiC.

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# Appendix A

## APPENDIX

### Derivation of effective carrier concentration

Resistance(R) of any material is given by the formula :

$$R = \frac{\rho L}{A} \dots (A.1)$$

where  $\rho$  is the resistivity of the material given as :

$$\rho = \frac{1}{Nq\mu} \dots (A.2)$$

### Over Region I

Referring to Fig.(4.1) considering the elemental horizontal strip. The resistance of the element over region I is given by eq.(4.8) is reproduced here again:

$$dR_I = \frac{dx}{A\mu_nq(N_o + \alpha_1x)} \dots (A.3)$$

Impurity profile(N(x)) over region I is given by eq.(4.6). Integrating above equation over region I.

$$R_I = \int_0^{h_1} \frac{dx}{A\mu_nq(N_o + \alpha_1x)} \dots (A.4)$$

the term  $A\mu_nq$  is constant (= k), eq.(A.4) reduces to:

$$R_I = \frac{1}{k} \int_0^{h_1} \frac{dx}{(N_o + \alpha_1 x)} \dots (A.5)$$

since

$$\int \frac{dx}{x} = \ln(x)$$

eq.(A.5) can be written as:

$$R_I = \frac{1}{\mu_n q A \alpha_1} \ln \left[ 1 + \frac{\alpha_1 h_1}{N_o} \right] \dots (A.6)$$

According to eqs.(A.1) and (A.2),  $R_I$  can also be written as:

$$R_I = \frac{h_1}{N_{eff1} q \mu_n A} \dots (A.7)$$

Equating eqs.(A.6) and (A.7) gives  $N_{eff1}$  as :

$$N_{eff1} = \frac{\alpha_1 h_1}{\ln \left[ 1 + \frac{\alpha_1 h_1}{N_o} \right]} \dots (A.8)$$

## Over Region II

Similar to region I, the resistance of the elemental strip over region II is given as :

$$dR_{II} = \frac{dx}{A\mu_nq(N_{o1} + \alpha_2 x)} \dots (A.9)$$

Impurity profile(N(x)) over region II is given by eq.(4.7). Integrating above equation over region II.

$$R_{II} = \int_{h_1}^{h_1+h_2} \frac{dx}{A\mu_nq(N_{o1} + \alpha_2 x)} \dots (A.10)$$

the term  $A\mu_nq$  is again considered constant = k, eq.(A.10) reduces to:

$$R_{II} = \frac{1}{k} \int_{h_1}^{h_1+h_2} \frac{dx}{(N_{o1} + \alpha_2x)} \dots (A.11)$$

eq.(A.11) can then be written as:

$$R_{II} = \frac{1}{\mu_nqA\alpha_2} \ln \left[ 1 + \frac{\alpha_2h_2}{(N_{o1} + \alpha_2h_1)} \right] \dots (A.12)$$

$R_{II}$  can also be written as(as per eqs.(A.1) and (A.2)):

$$R_{II} = \frac{h_2}{N_{eff2}q\mu_nA} \dots (A.13)$$

Equating eqs.(A.12) and (A.13) gives  $N_{eff2}$  as :

$$N_{eff2} = \frac{\alpha_2h_2}{\ln \left[ 1 + \frac{\alpha_2h_2}{(N_{o1} + \alpha_2h_1)} \right]} \dots (A.14)$$

Hence, the total effective carrier concentration is given as summation of eqs.(A.8) and (A.14):

$$N_{eff} = \frac{\alpha_1h_1}{\ln \left[ 1 + \frac{\alpha_1h_1}{N_o} \right]} + \frac{\alpha_2h_2}{\ln \left[ 1 + \frac{\alpha_2h_2}{N_{o1} + \alpha_2h_1} \right]} \dots (A.15)$$

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